

## **MTH6132: Relativity**

**Duration: 2 hours**

**Date and time: 5th May 2016, 10:00–12:00**

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**Apart from this page, you are not permitted to read the contents of this question paper until instructed to do so by an invigilator.**

**You should attempt ALL questions. Marks awarded are shown next to the questions.**

**Calculators are not permitted in this examination. The unauthorised use of a calculator constitutes an examination offence.**

Complete all rough workings in the answer book and cross through any work that is not to be assessed.

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**Exam papers must not be removed from the examination room.**

**Examiner(s): S. Beheshti**

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**Question 1.** Let  $F$  and  $F'$  denote two inertial reference systems moving with velocity  $v$  with respect to each other. In  $F$ , two events occur simultaneously at  $t = 0$ , separated by a distance  $X$  along the  $x$ -axis. The time interval between the events in  $F'$  is  $T$ .

- (a) Draw a 2-dimensional spacetime diagram describing the situation, including both  $F$  and  $F'$ . You may assume units for which  $c = 1$ . [4]
- (b) Show that the spatial distance between the two events in  $F'$  is  $\sqrt{X^2 + T^2}$ . [4]
- (c) Determine the relative velocity  $v$  of the frames  $F$ ,  $F'$  in terms of  $X$  and  $T$ . You may assume  $c = 1$  in your calculations. [7]

**Question 2.** Let  $\bar{A}$  and  $\bar{B}$  denote two arbitrary 4-vectors in Minkowski spacetime.

- (a) Define what is meant by the scalar product  $\bar{A} \cdot \bar{B}$ . What does it mean to say  $|\bar{A}|^2$  is an invariant? [4]
- (b) Using the fact that  $|\bar{A}|^2$ ,  $|\bar{B}|^2$  and  $|\bar{A} + \bar{B}|^2$  are invariants, show that the scalar product  $\bar{A} \cdot \bar{B}$  is also an invariant. [4]
- (c) Show that the sum of any two orthogonal spacelike vectors is also spacelike. [7]

**Question 3.** The metric for a particular two-dimensional spacetime is given by

$$ds^2 = \frac{1}{y^2} (dx^2 + dy^2)$$

- (a) Determine the covariant and contravariant components of the metric tensor for this spacetime. [2]
- (b) Employ the formula for the Christoffel symbols given in the Appendix to calculate the components  $\Gamma^2_{11}$ ,  $\Gamma^1_{12}$  and  $\Gamma^2_{22}$  of the connection for this metric. [Note the identification  $(x^1, x^2) = (x, y)$  is used here.] [6]
- (c) Assuming that the components in part (b) are the only nonzero ones, confirm that the  $R^2_{121}$  component of the Riemann tensor for this metric is  $-\frac{1}{y^2}$ . Given that Gauss curvature is given by  $K = R_{1212}/(g_{11}g_{22} - g_{12}g_{21})$ , can this metric describe a flat spacetime? [7]

**Question 4.**

- (a) Assume  $\nabla$  is the Levi-Civita connection of a metric  $g_{ab}$ . Using properties of this covariant derivative, simplify fully the expression

$$\nabla_a (g_{bc} S^{bc}).$$

[3]

- (b) Let  $X^a$  be the tangent vector to a geodesic given by  $x^a(\lambda)$ . Using part (a), show that the norm of this tangent vector is conserved along geodesics, i.e.,

$$X^a \nabla_a (|X|^2) = 0.$$

[7]

**Question 5.** In this question consider units for which  $c = 1$ . A particle has rest mass  $m_0$ . Whilst at rest, it emits a photon and, as a result, its rest mass is reduced to  $m_0/2$ . By comparing components of the 4-momenta before and after the event, show that the speed of the particle after the reduction of mass is  $3/5$ . Show also that the energy  $E = h\nu$  of the photon is  $3m_0/8$ .

[10]

**Question 6.**

- (a) It can be shown that in a Local Inertial Frame the Riemann tensor can be expressed in the form

$$R_{abcd} = \frac{1}{2} (\partial_d \partial_a g_{bc} + \partial_c \partial_b g_{ad} - \partial_c \partial_a g_{bd} - \partial_d \partial_b g_{ac})$$

at a specific point  $p$ . Explain what is meant by a Local Inertial Frame and employ the expression above to show that

$$R_{abcd} = -R_{bacd}$$

at  $p$ . Is this relation valid in an arbitrary frame of reference? Explain your reasoning.

[5]

- (b) Suppose that the curvature of a spacetime satisfies the equation

$$R_{ab} - \frac{1}{2} R g_{ab} + \lambda g_{ab} = 0,$$

where  $\lambda$  is a constant. Define the Ricci scalar and show that it satisfies

$$R = 4\lambda.$$

Hint: You will also need to prove and use the fact that  $g_{ab} g^{ab} = 4$ .

[5]

**Question 7.**

- (a) Write down the transformation laws under general coordinate transformations for a  $(1, 0)$ -tensor and a  $(0, 2)$ -tensor, respectively. Use this to show that the product of these tensors is a tensor of type  $(1, 2)$ . [5]
- (b) Prove that if  $W_{ab}{}^c$  is a  $(1, 2)$ -tensor, then  $W_{ab}{}^b$  is a  $(0, 1)$ -tensor. [5]

**Question 8.** Consider the Schwarzschild metric

$$ds^2 = - \left( 1 - \frac{2GM}{r} \right) dt^2 + \left( 1 - \frac{2GM}{r} \right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2).$$

- (a) What physical situation is described by this metric? What happens if  $M = 0$ ? [2]
- (b) At what two values of  $r$  is this metric in the above form singular? Re-express the Schwarzschild metric in terms of *Eddington-Finkelstein coordinates*, given by  $(\hat{t}, r, \theta, \varphi)$ , where

$$\hat{t} = t + 2GM \ln |r - 2GM|.$$

Show that in these coordinates, one of the two singularities is removed.

- [6]
- (c) Use the Euler-Lagrange equations to derive the geodesic equations obeyed by a photon for this metric. [7]

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**End of Paper—An appendix of 1 page follows.**

You are reminded of the following information, which you may use without proof.

- Lower case Latin indices run from 0 to 3.
- The metric tensor of the Minkowski spacetime is  $\eta_{ab}$  such that

$$ds^2 = \eta_{ab} dx^a dx^b = -c^2 dt^2 + dx^2 + dy^2 + dz^2$$

- The Lorentz transformations between two frames  $F$  and  $F'$  in standard configuration are given by

$$x' = \gamma(x - vt), \quad t' = \gamma \left( t - \frac{vx}{c^2} \right), \quad y' = y, \quad z' = z$$

where

$$\gamma = \frac{1}{\sqrt{1 - (v^2/c^2)}}$$

and  $F'$  is moving with speed  $v$  relative to  $F$ .

- The covariant derivative of a covariant vector is given by

$$\nabla_a V_b = \partial_a V_b - \Gamma^f_{ba} V_f.$$

- The covariant derivative of a contravariant vector is given by

$$\nabla_a V^b = \partial_a V^b + \Gamma^b_{af} V^f.$$

- The metric tensor satisfies:

$$g_{ab} g^{bc} = \delta_a^c.$$

- Christoffel symbols (connection):

$$\Gamma^m_{ij} = \frac{1}{2} g^{mk} (\partial_i g_{kj} + \partial_j g_{ik} - \partial_k g_{ij}).$$

- The Riemann curvature tensor:

$$R^a_{bcd} = \partial_c \Gamma^a_{bd} - \partial_d \Gamma^a_{bc} + \Gamma^a_{ec} \Gamma^e_{bd} - \Gamma^a_{ed} \Gamma^e_{bc}.$$

- Euler–Lagrange equations:

$$\frac{d}{d\lambda} \left( \frac{\partial L}{\partial \dot{x}^c} \right) - \frac{\partial L}{\partial x^c} = 0$$

- Geodesic equations:

$$\ddot{x}^a + \Gamma^a_{bc} \dot{x}^b \dot{x}^c = 0.$$

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**End of Appendix.**