Diffusion coefficient

- The net flux from left to right therefore is: $J_Z = J(L \to R) J(L \leftarrow R) = \frac{1}{4} v_{mean} \{ \mathcal{N}(-\lambda) \mathcal{N}(\lambda) \} = \frac{1}{4} v_{mean} \left\{ \left[\mathcal{N}(0) \lambda \left(\frac{d\mathcal{N}}{dz} \right)_0 \right] \left[\mathcal{N}(0) + \lambda \left(\frac{d\mathcal{N}}{dz} \right)_0 \right] \right\} = -\frac{1}{2} v_{mean} \lambda \left(\frac{d\mathcal{N}}{dz} \right)_0$, in accord with Fick's 1st law.
- From this we should be able pick out the diffusion coefficient as: $\frac{1}{2}v_{mean}\lambda$, however, even if a MQ begins its journey near the window it could have a long flight before it gets there so in reality we have to multiply by a factor of 2/3, so:
- $D = \frac{1}{3} \lambda v_{mean}$, diffusion coefficient.

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Thermal conductivity (reminder)



- According to equipartition theorem, each MQ carried an average energy ε=vkT, where v is a number of the
 order of 1. For atoms, v=3/2.
- When a MQ passes through the imaginary window, it transports that average energy. We assume uniform $\mathcal N$ and a T gradient.
- MQs arriving from the left travel a λ from their last collision in a hotter region, and so with higher energy.
 MQs also arrive from the right after travelling a λ from a cooler region.
- The 2 opposing energy fluxes are: $J(L \to R) = \frac{1}{4} \mathcal{N}v_{mean} \varepsilon(-\lambda)$ and $J(L \leftarrow R) = \frac{1}{4} \mathcal{N}v_{mean} \varepsilon(\lambda)$, with Z_W
- v is a number of the order of 1; for atoms it's 3/2.