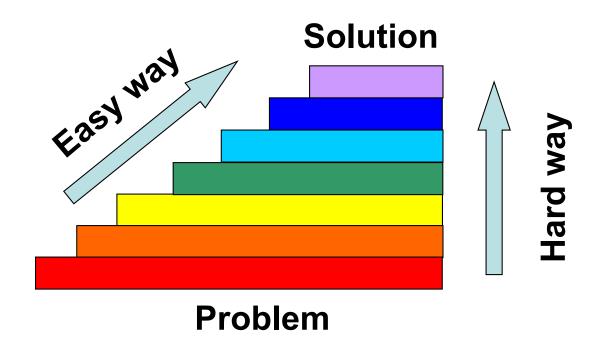
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Renewable Energy Sources

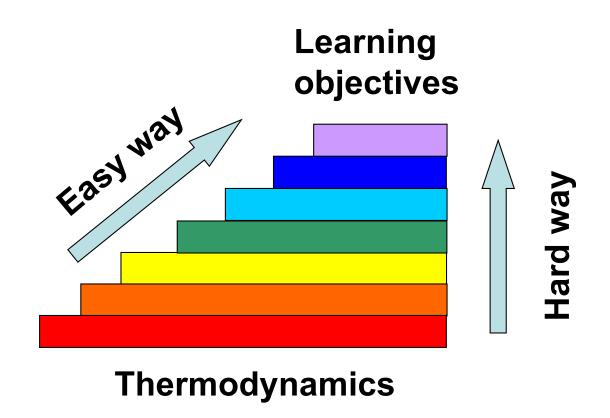
Thermodynamics of Energy Conversion - REVISION

Prof. Huasheng Wang

A step-by-step approach



A step-by-step approach



- General step-by-step approach
- 1. Problem statement; info given, to be found, objectives
- 2. Draw a schematic; indicate relevant info on
- 3. Assumptions and approximations; simplify, justify
- 4. Physical laws; apply relevant laws and principles
- 5. Properties; property relations, tables, property source
- 6. Calculations; units, significant number ("do not copy")
- 7. Reasoning, verification, and discussion; validity, significance, implications, conclusions, recommendations, limitations,
- 8. Neatness, organization, completeness, visual appearance

 Conservation of mass, momentum, energy & electrical charge (the first law of thermodynamics)

$$\begin{cases}
\text{The net gain in} \\
X \text{ by the system}
\end{cases} = \begin{cases}
\text{The net amount of } X \\
\text{transported into the system}
\end{cases}$$

$$+ \begin{cases}
\text{The net amount of } X \\
\text{generated by the sytem}
\end{cases}$$

$$1+2=3$$
 $a+b=c$
 $1+2-3=0$ $a+b-c=0$
 $1=3-2$ $a=c-b$

Balance of entropy (ΔS_{irr})

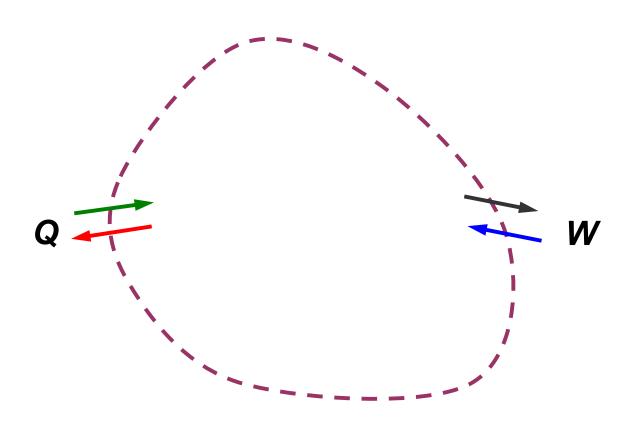
$$\begin{cases}
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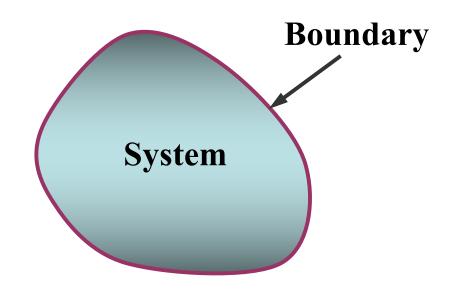
$$a + b = c$$

 $a + b - c = 0$
 $a = c - b$

• Sign "convention"



Thermodynamic system
material, body, amount of fluid, under discussion,
identifiable collection of matter



Surroundings/Environments

Thermodynamic system

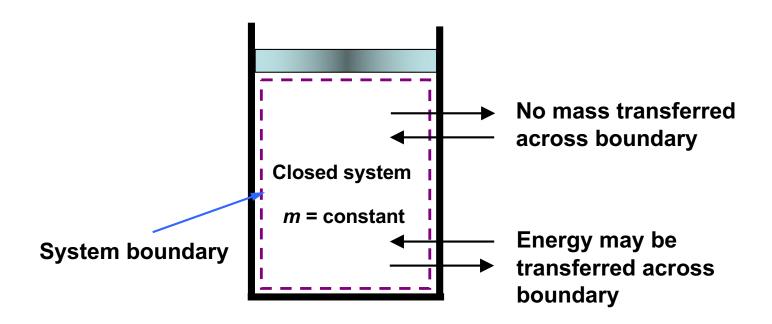
Surroundings/Environments

Boundary

- Fixed and moving
- Real and imaginary; work
- "Thermal conductor" and adiabatic; heat
- Permeable and non-permeable; mass

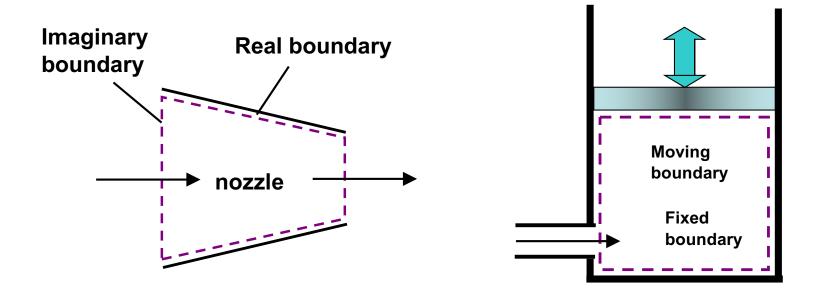
Thermodynamic system

Closed system (control mass)



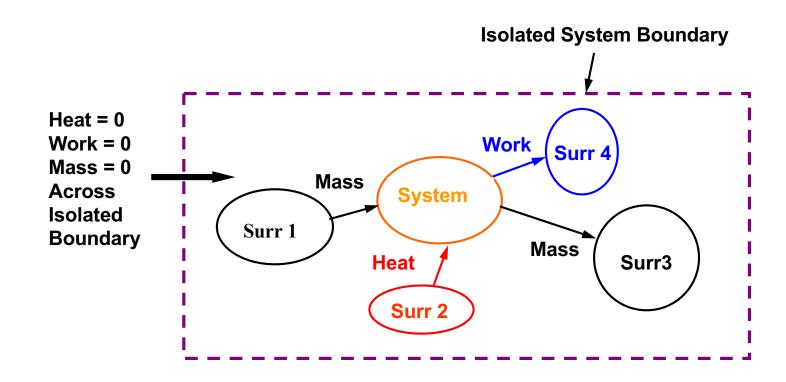
Thermodynamic system

Open system (control volume)



Thermodynamic system

Isolated system



Property of a system
 Characteristic of a system which can be measured

Extensive properties
 mass m, volume V, energy E, entropy S, etc.

Intensive properties
 pressure *P*, temperature *T*, specific volume *v*, etc.

Energy, E and entropy, S

•
$$dE = 0$$

• dS≥0

For an isolated system the internal energy is constant and the entropy can only increase. Entropy is constant for ideal reversible process

Property of a system

E, S, V fundamental thermodynamic properties (extensive)

Specific quantity – quantity / mass of system

Specific properties

$$v = V/m$$
, $e = E/m$, $s = S/m$

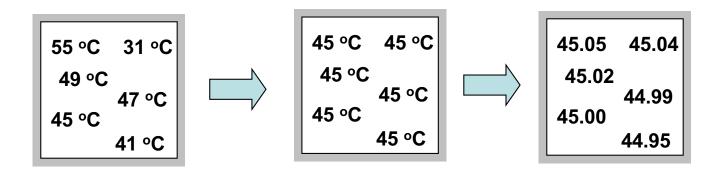
Other specific quantities

$$w = W/m$$
, $q = Q/m$

Specific force of gravity mg/m = g

Equilibrium

The properties of systems ("left alone" i.e. isolated) are observed to attain constant values, i.e. the system comes to equilibrium. Subsystems in equilibrium are themselves in equilibrium



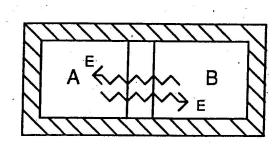
Equilibrium state

When all of the properties of a system are constant

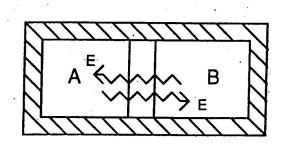
How many independent properties are needed in order to determine a equilibrium state?

- Simple fluid
- $\psi(E, S, V) = 0$
- S = S(E, V)
- E = E(S, V)
- For unit mass of a simple homogenous fluid
- $\Phi(e, s, v) = 0$
- s = s(e, v)
- e = e(s, v)

- Temperature and pressure
- Consider isolated system shown comprising two subsystems A and B
- A + B isolated; A and B not isolated from each other
- Entropy is extensive so $S = S_A + A_B$
- System is isolated so
- $E_A + E_B = constant$
- $V_A + V_B = constant$



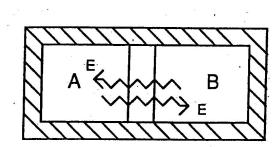
- Temperature and pressure
- (a) Energy may be shared between A and B if the dividing boundary permits energy transfer (thermal conductor).
- (b) Volume may be shared between A and B if the dividing boundary is freely movable.
- The equilibrium state of the isolated system must be an equilibrium state of A and an equilibrium state of B with energy and volume distributed between the two so as to maximize the total entropy



- Temperature and pressure
- It may be shown that when (a) is true

$$(\partial E/\partial S)_V = (\partial E/\partial S)_V$$

System A System B

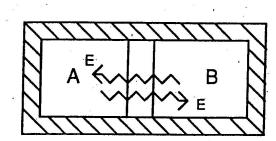


 Thermodynamic temperature of a simple fluid system is defined by

$$T = (\partial E/\partial S)_V = (\partial e/\partial s)_V$$

- Temperature and pressure
- It may be shown that when (b) is true

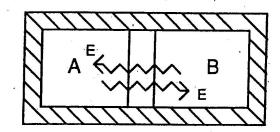
$$(\partial E/\partial V)_S$$
 = $(\partial E/\partial V)_S$
System A System B



Thermodynamic pressure of a simple fluid system is defined by

$$P = -(\partial E/\partial V)_S = -(\partial e/\partial v)_S$$

- Temperature and pressure
- So, for equilibrium

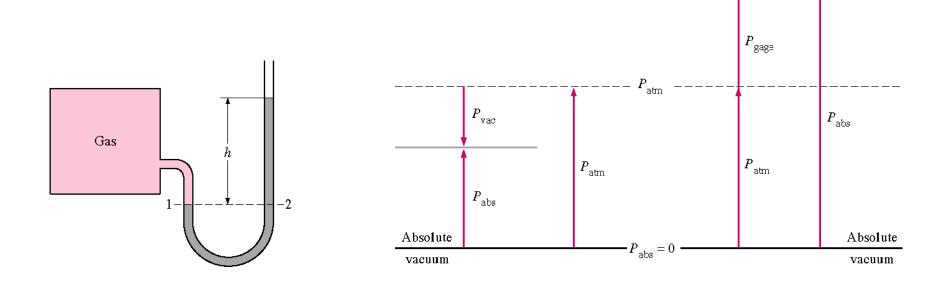


- $T_A = T_B$ when separating boundary is a "thermal conductor"
- $P_A = P_B$ when separating boundary is freely moveable
- Driving force

$$E = E(S, V)$$
 $dE = \left(\frac{\partial E}{\partial S}\right)_{V} dS + \left(\frac{\partial E}{\partial V}\right)_{S} dV$

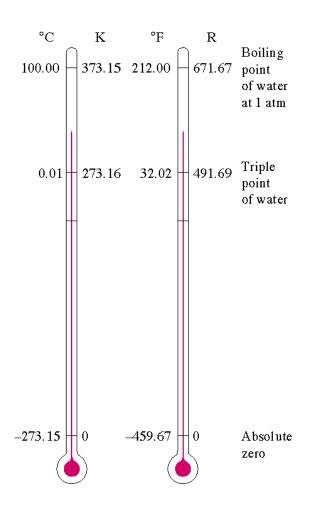
$$dE = TdS - PdV$$
 $de = Tds - Pdv$

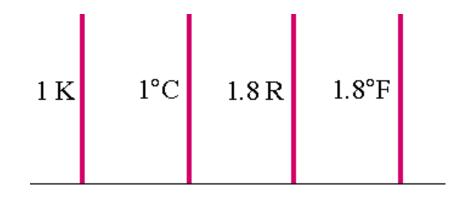
Pressure



• $P_{abs} = P_{atm} + P_{gauge}$ (Be careful!)

Temperature





- T = t + 273.15
- Absolute temperature should be used in all thermodynamics calculation.

Two-property rule for closed simple fluid system

For equilibrium states

•
$$E = E(S, V)$$
 (1)

•
$$T = (\partial E/\partial S)_V = T(S, V)$$
 (2)

•
$$P = -(\partial E/\partial V)_S = P(S, V)$$
 (3)

Homogeneous closed system

$$\bullet \quad \mathbf{e} = \mathbf{e}(\mathbf{s}, \, \mathbf{v}) \tag{1}$$

•
$$T = (\partial e/\partial s)_v = T(s, v)$$
 (2)

•
$$P = -(\partial e/\partial v)_s = P(s, v)$$
 (3)

Gibbs Phase Rule

- Multicomponent, multiphse system
- IV = C PH + 2
- IV the number of independent variables
- C the number of components (constituents, species)
- PH the number of phases present in equilibrium

Summary – Property rule

Closed system (constant mass)

State is fixed by any 2 of E, S, V, P, T, G, H

- Homogeneous closed system
 Intensive state is fixed by any 2 of e, s, v, P, T, g, h
- Homogeneous phase of pure fluid
 Intensive state is fixed by any 2 of e, s, v, P, T, g, h

Specific Heat-Capacities

Summary

$$\boldsymbol{c}_{v} = \boldsymbol{T} \left(\frac{\partial \mathbf{s}}{\partial \boldsymbol{T}} \right)_{v} = \left(\frac{\partial \mathbf{e}}{\partial \boldsymbol{T}} \right)_{v} \qquad \boldsymbol{c}_{P} = \boldsymbol{T} \left(\frac{\partial \mathbf{s}}{\partial \boldsymbol{T}} \right)_{P} = \left(\frac{\partial \boldsymbol{h}}{\partial \boldsymbol{T}} \right)_{P}$$

• If c_v can be taken constant between T_1 and T_2 , then, for a quasi-equilibrium constant volume process

$$q = c_v (T_2 - T_1) = e_2 - e_1$$

• If c_p can be taken constant between T_1 and T_2 , then, for a quasi-equilibrium constant pressure process

$$q = c_P(T_2 - T_1) = h_2 - h_1$$

Ideal Gas - 1

- For real gases at low pressure Pv/T≈ constant
- · We define an ideal gas as a fluid for which

$$Pv/T = R = constant$$

- Therefore real gases at low pressures approximate to ideal gases. R has a different value for different gases and is called the specific ideal-gas constant
- Ideal gas (Chemistry) a hypothetical gas whose molecules occupy negligible space and have no interactions, and which consequently obeys the gas law exactly.

Ideal Gas - 2

Summary

$$Pv/T = R \text{ (constant)} \qquad P_1v_1/T_1 = P_2v_2/T_2$$

$$e = e(T) \qquad h = h(T)$$

$$c_v = \frac{de}{dT} \qquad c_P = \frac{dh}{dT} \qquad c_P - c_v = R$$

$$e_2 - e_1 = \int_1^2 c_v dT \qquad h_2 - h_1 = \int_1^2 c_P dT$$

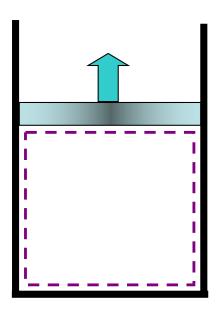
$$s_2 - s_1 = \int_1^2 \frac{de}{T} + R \ln \left(\frac{v_2}{v_1} \right) = \int_1^2 \frac{dh}{T} - R \ln \left(\frac{P_2}{P_1} \right)$$

Perfect Gas

- Summary
- Pv = RT and c_v , c_P constant
- $e_2 e_1 = c_v(T_2 T_1)$
- $h_2 h_1 = c_P(T_2 T_1)$
- $s_2 s_1 = 3$ expressions in terms of (T, v), (T, P), (P, v)
- For isentropic (constant s) process
- Pv^{γ} = constant; $P_1v_1^{\gamma} = P_2v_2^{\gamma}$
- $\gamma = c_P/c_V = \text{constant for perfect gas}$

Quasi-equilibrium (quasi static, non-dissipative, internally reversible) process

Ideal process performed so slowly that system is always in an equilibrium state i.e. during the process the system passes through a continuous sequence of equilibrium states, e.g. compression of a gas in a cylinder *infinitely slowly*.



Summary

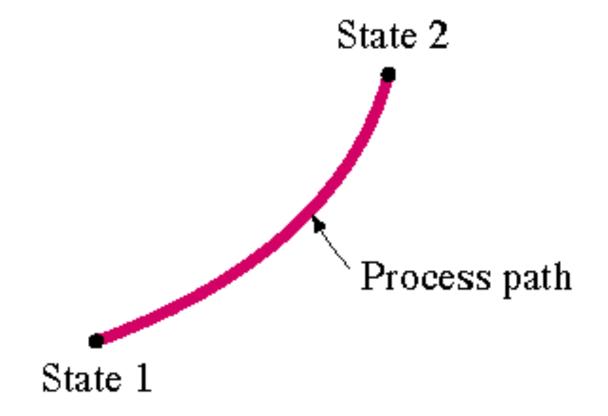
$$E_2 - E_1 = Q + W$$
 always true

$$W = -\int_{V_1}^{V_2} P dV$$
 finite quasi-equilibrium process

$$Q = \int_{S_1}^{S_2} T dS$$
 finite quasi-equilibrium process

Process

Change from one equilibrium state to another

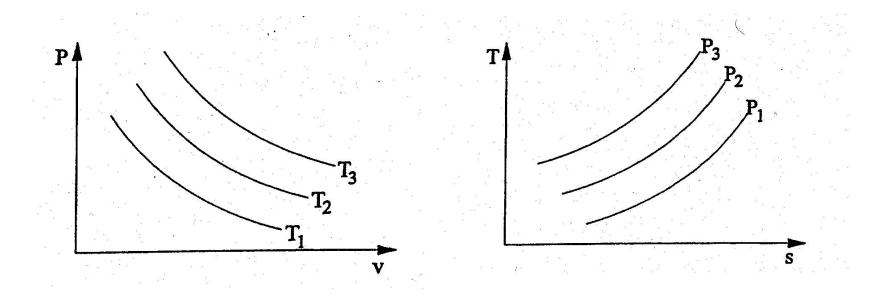


States between 2 end states can be equilibrium or non-equilibrium.

Work and Heat later.

Process illustration

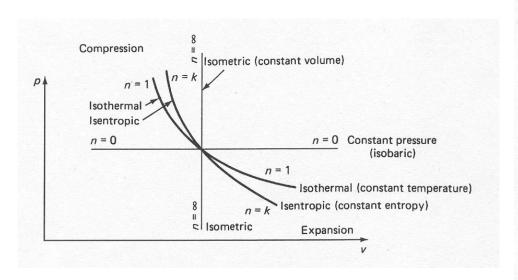
P-V and T-S diagrams

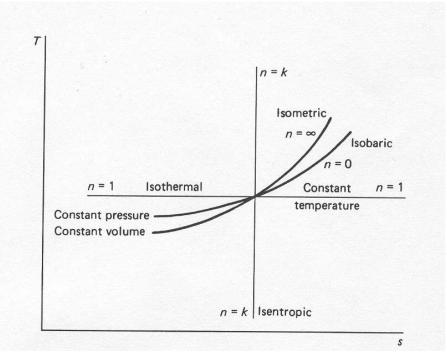


Polytropic process

```
    Pv<sup>n</sup> = constant -∞< n <∞</li>
    Isobaric P = constant 0
    Isothermal T = constant 1
    Isentropic S = constant γ
    Isochoric V = constant -/+∞
```

Polytropic Process





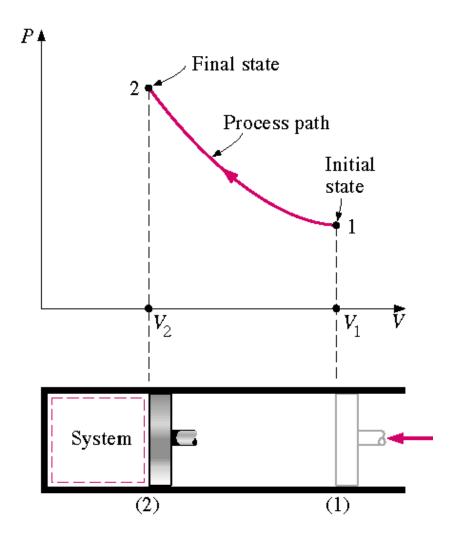
Polytropic Process

Table 1 Summary of perfect gas relations

Process Isochoric v=constant $P=constant$ Isothermal $P=constant$ $P=$		abic i ouiiiii	ially of periec	i gas relations		
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	Process			T = constant	s = constant $Pv^{\gamma} = const.$ $Tv^{\gamma-1} = constant$	Polytropic $Pv^{n} = \text{constant}$ $Tv^{n-1} = \text{constant}$ $TP^{-\frac{n-1}{n}} = \text{constant}$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	n	∞	0	1		-∞ to +∞
relation $\frac{T_2}{T_2} = \frac{T_2}{P_2}$ $\frac{T_2}{T_2} = \frac{V_2}{V_2}$ $\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)$ $\frac{T_2}{P_1} = \left(\frac{V_1}{V_2}\right)$ $\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)$ $\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)$ $\frac{T_2}{T_1} = \left(\frac{P_2}{P_1}\right)^{\frac{\gamma-1}{\gamma}}$ $\frac{T_2}{T_1} = \left(\frac{P_2}{P_1}\right)$ $\frac{U_2 - U_1}{P_1} = \frac{C_v(T_2 - T_1)}{P_1}$ $\frac{C_v(T_2 - T_1)}{P_2} = \frac{C_v(T_2 - T_1)}{P_1}$ $\frac{C_v(T_2 - T_1)}{P_2} = \frac{C_v(T_2 - T_1)}{P_2}$		Cv				$c_n = c_v \left(\frac{\gamma - n}{1 - n} \right)$
$\frac{T_{2}}{T_{1}} = \left(\frac{P_{2}}{P_{1}}\right)^{\frac{\gamma-1}{\gamma}} \qquad \frac{T_{2}}{T_{1}} = \left(\frac{P_{2}}{P_{1}}\right)^{\frac{\gamma-1}{\gamma}} \qquad \frac{C_{\nu}(T_{2}-T_{1})}{C_{\nu}(T_{2}-T_{1})} \qquad \frac{C_{\nu}(T_{2}-T_$		$\frac{T_1}{T_2} = \frac{P_1}{P_2}$	$\frac{T_1}{T_2} = \frac{V_1}{V_2}$	$P_1V_1=P_2V_2$	$\frac{P_2}{P_1} = \left(\frac{v_1}{v_2}\right)^r$	$\frac{P_2}{P_1} = \left(\frac{v_1}{v_2}\right)^n$
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$					1 (2)	$\frac{T_2}{T_1} = \left(\frac{v_1}{v_2}\right)^{n-1}$
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$					$\frac{T_2}{T_1} = \left(\frac{P_2}{P_1}\right)^{\frac{\gamma-1}{\gamma}}$	$\frac{T_2}{T_1} = \left(\frac{P_2}{P_1}\right)^{\frac{n-1}{n}}$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	u ₂ – u ₁	$c_v(T_2-T_1)$	$c_v(T_2-T_1)$	0	$c_v(T_2-T_1)$	$c_v(T_2 - T_1)$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$h_2 - h_1$	$c_p(T_2-T_1)$	$c_p(T_2-T_1)$	0	$c_p(T_2-T_1)$	$c_p(T_2-T_1)$
$r_1v_1mv_1$ $1-k$ $1-n$	s ₂ - s ₁	$c_v \ln \frac{T_2}{T_1}$	$c_p \ln \frac{T_2}{T_1}$	$R \ln \frac{v_2}{v_1}$	0	$c_n \ln \frac{T_2}{T_1}$
$q = c_v(T_2 - T_1) = c_p(T_2 - T_1) = P_1 v_1 \ln \frac{v_2}{v_1} = 0$ $c_n(T_2 - T_2) = 0$	W	0	$P(v_2 - v_1)$	$P_1 v_1 \ln \frac{v_2}{v_1}$	$\frac{P_2V_2 - P_1V_1}{1 - k}$	$\frac{P_2V_2 - P_1V_1}{1-n}$
	q	$c_{v}(T_{2}-T_{1})$	$c_p(T_2-T_1)$	$P_1 v_1 \ln \frac{v_2}{v_1}$	0	$c_n(T_2-T_1)$

Process

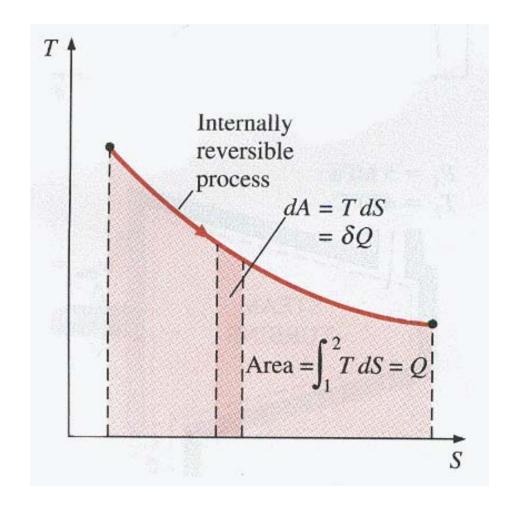
Process illustration (P – V diagram)



Property?

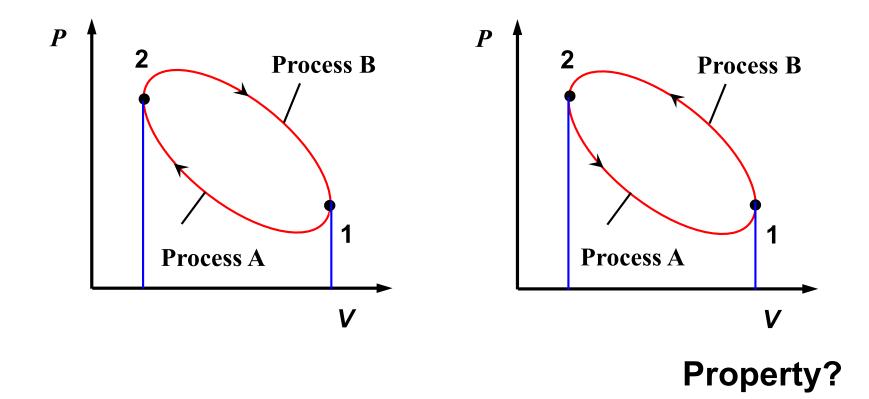
Process

Process illustration (T – S diagram)

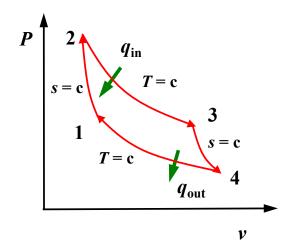


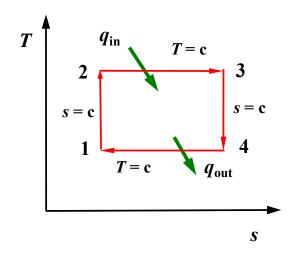
Cycle

A serious of a connected processes with identical end states



Carnot cycle – an example



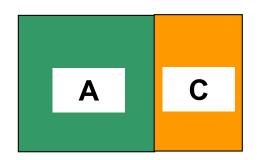


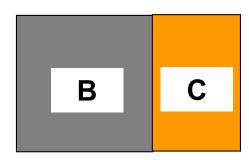
Thermal efficiency = Desired output/Required input

COP (Coefficient Of Performance) = Desired output/Required input

Laws of thermodynamics - 1

The zeroth law of thermodynamics
 R. H. Fowler (1931)





Laws of thermodynamics - 2

The first law of thermodynamics

For any process of a stationary system

$$W + Q = E_2 - E_1$$

Steady flow energy equation (SFEE)

$$W + Q = \Sigma [m(h + v^2/2 + gz)]_2 - \Sigma [m(h + v^2/2 + gz)]_1$$

Laws of thermodynamics - 3

Entropy and the 2nd Law of thermodynamics

Spontaneous (no outside influence) processes lead to increase in entropy. This is the essence of the Second Law of Thermodynamics.

• dE = 0

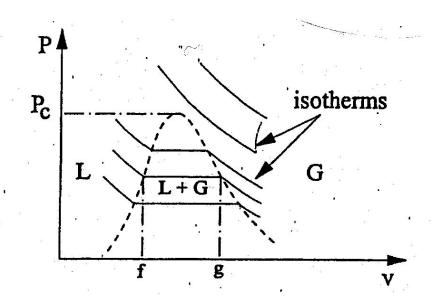
• dS≥0

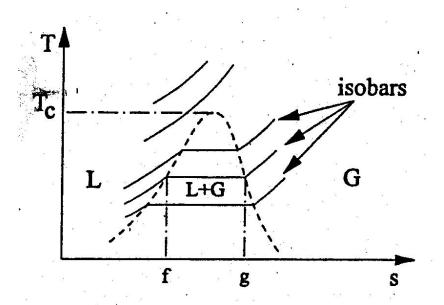
For an <u>isolated system</u> the internal energy is constant and the entropy can only increase. Entropy is constant for ideal

reversible process

Property of Pure Substances

• The P-V-T Surface for a Real Substance





Average specific properties for 2-phase states

"Quality" or "dryness" x is defined

so that
$$y = \frac{m_f y_f + m_g y_g}{m_f + m_g}$$

= $(1-x)y_f + xy_g$

Note: at
$$x = 0$$
, $y = y_f$
 $x = 1$, $y = y_g$

We also define, for convenience,

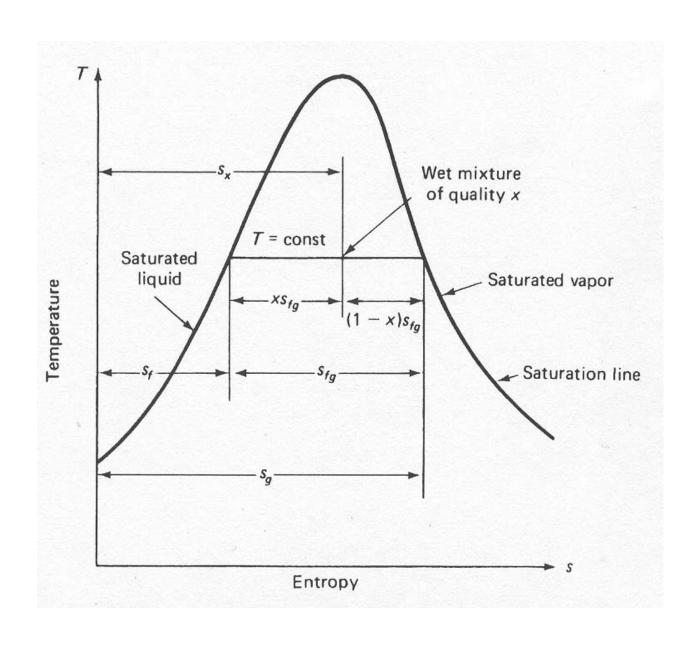
$$oldsymbol{y}_{\mathsf{fg}} = oldsymbol{y}_{\mathsf{g}} - oldsymbol{y}_{\mathsf{f}}$$

so that

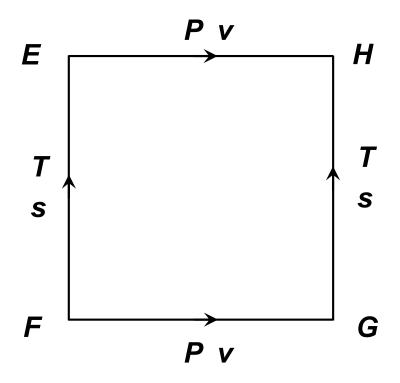
$$y = (1-x)y_f + x(y_{fg} + y_f)$$

= $y_f + xy_{fg}$

Average specific properties for 2-phase states

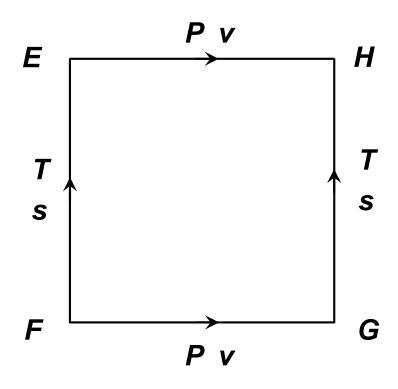


Relation of properties - 1



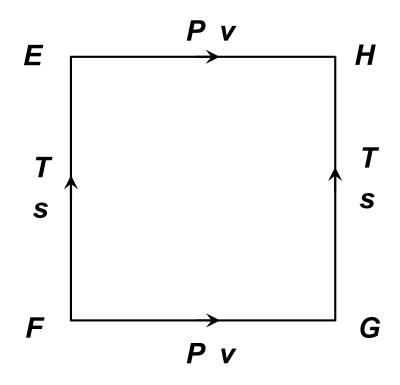
F – Helmholtz function; G – Gibbs function; H – enthalpy

Relation of properties - 2



$$E = F + TS$$
 $E = H - PV$
 $F = E - TS$ $F = G - PV$
 $G = F + PV$ $G = H - TS$
 $H = E + PV$ $H = G + TS$

Relation of properties - 3



$$de = Tds - Pdv$$

$$dh = Tds + vdP$$

$$df = -sdT - Pdv$$

$$dg = -sdT + vdP$$

$$d(corner) = \sum (near)d(far)$$

References

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