

① Selected solutions to problem set 6

2. The solution is

$$\begin{aligned}
 u(x,t) &= \frac{1}{2} [e^{-(x-ct)^2} + e^{-(x+ct)^2}] \\
 &\quad + \frac{1}{2c} \int_{x-ct}^{x+ct} 1 \, ds \\
 &\quad + \frac{1}{2c} \int_0^t \int_{x-ct+cs}^{x+ct+cs} (-r) \, dr \, ds \\
 &= \frac{1}{2} [e^{-(x-ct)^2} + e^{-(x+ct)^2}] + t \\
 &\quad - \cancel{\frac{1}{2c} \int_0^t \frac{(x+ct+cs)^2 - (x-ct+cs)^2}{2} \, ds} \\
 &= \frac{1}{2} [e^{-(x-ct)^2} + e^{-(x+ct)^2}] + t \\
 &\quad + \frac{(x+ct+cs)^3}{12c^2} \Big|_0^t + \frac{(x-ct+cs)^3}{12c^2} \Big|_0^t \\
 &= \frac{1}{2} [e^{-(x-ct)^2} + e^{-(x+ct)^2}] + t \\
 &\quad + \frac{x^3}{6c^2} + \frac{(x+ct)^3 + (x-ct)^3}{12c^2}
 \end{aligned}$$

②

3. The solution is

$$u(x,t) = \frac{1}{2c} \int_0^t \int_{x-ct+cs}^{x+ct-cs} r dr ds$$

$$= \frac{1}{2c} \int_0^t [(x+ct-cs) - (x-ct+cs)] ds$$

$$= \frac{1}{2c} \int_0^t (2ct - 2cs) ds$$

$$= \frac{1}{2c} \cdot [2cts - cs^2] \Big|_0^t$$

$$= \frac{1}{2c} [2ct^2 - ct^2]$$

$$= \frac{t^2}{2}$$

③ 5. The general solutions are.

$$u(x,t) = \sum_{n=1}^{\infty} a_n \sin(nx) \cos(nt) + b_n \sin(nx) \sin(nt)$$

with

$$a_n = \frac{2}{\pi} \int_0^\pi \sin x \cdot \sin(nx) dx$$

$$= \begin{cases} 0 & , n \neq 1 \\ \frac{2}{\pi} \int_0^\pi \frac{1 - \cos(2nx)}{2} dx & , n=1 \end{cases}$$

$$= \begin{cases} 0 & , n \neq 1 \\ 1 & , n=1 \end{cases}$$

So the solution is

$$u(x,t) = \sin x \cos(ct)$$

6. (1) For $x > 0$,

$$F(x) = f(x)$$

and because $-x < 0$,

$$F(-x) = f(-(-x)) = f(x) = F(x)$$

so F is even.

G is even by the same reason.

~~The derivative of an even function is always odd.~~

(2) * The solution is

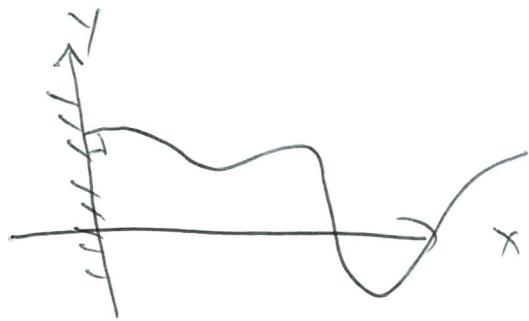
$$\begin{aligned} V(x, t) &= \frac{1}{2} [F(x+ct) + F(x-ct)] \\ &\quad + \frac{1}{2c} \int_{x-ct}^{x+ct} G(s) ds \end{aligned}$$

$$\begin{aligned} (3) \quad V_x(x, t) &= \frac{1}{2} [F'(x+ct) + F'(x-ct)] \\ &\quad + \frac{1}{2c} G(x+ct) - \frac{1}{2c} G(x-ct) \end{aligned}$$

$$\begin{aligned} (4) \quad V_x(0, t) &= \frac{1}{2} [F'(ct) + F'(-ct)] \\ &\quad + \frac{1}{2c} [G(ct) - G(-ct)] \\ &= 0 \end{aligned}$$

because F' is odd and G is even.

(5)

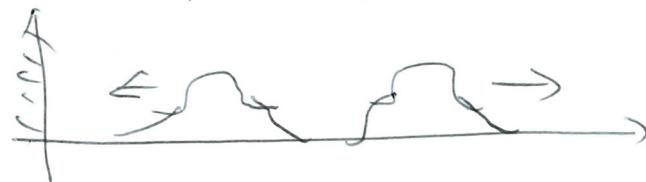


The string is allowed to move up and down, but always orthogonal to the y axis at the left end.

For an initial bump



it first separate into 2 bumps



~~at~~

After some time the left bump hit the "wall" (y -axis) and is reflected. Then both bumps are propagating to right



④ 8. . The Energy is

$$E[u](t) = \int_{-\infty}^{\infty} \left[\frac{1}{2} u_t^2 + \frac{1}{2} a^2 u_x^2 \right] dx$$

Multiply both sides of equation by u_t , we get

$$u_t \cdot u_{tt} - a^2 u_t u_{xx} - c u_t^2 = 0$$

Integrate get

$$\int_{-\infty}^{\infty} u_t \cdot u_{tt} dx - \int_{-\infty}^{\infty} a^2 u_t u_{xx} dx - c^2 \int_{-\infty}^{\infty} u_t^2 dx = 0$$

$$\cancel{\int_{-\infty}^{\infty} u_t \cdot u_{tt} dx} - \cancel{\int_{-\infty}^{\infty} a^2 u_t u_{xx} dx} - \int_{-\infty}^{\infty} a^2 u_t u_{xx} - c^2 \int_{-\infty}^{\infty} u_t^2 dx = 0$$

By integration by part., we get

$$\int_{-\infty}^{\infty} \frac{(u_t^2)_x}{2} dx - a^2 u_t u_x \Big|_{-\infty}^{\infty} + \int_{-\infty}^{\infty} a^2 u_{xx} u_x dx - c^2 \int_{-\infty}^{\infty} u_t^2 dx = 0$$

thus

$$(5) \frac{d}{dt} \int_{-\infty}^{\infty} \frac{(u_t)^2}{2} dx - 0 + \frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} \frac{a^2(u_x)^2}{2} dx - c \int_{-\infty}^{\infty} u_t^2 = 0.$$

Namely

$$\frac{d}{dt} \int_{-\infty}^{\infty} \left(\frac{u_t^2}{2} + \frac{a^2 u_x^2}{2} \right) dx = c \int_{-\infty}^{\infty} u_t^2 \leq 0$$

so the energy is non-increasing

because $c < 0$.

• Suppose u_1, u_2 are two solutions
to $\begin{cases} u_{tt} - a^2 u_{xx} - c u_t = f(x), \\ u(x, 0) = g(x), \\ u_t(x, 0) = h(x) \end{cases}$

then $u = u_1 - u_2$ is a solution to

$$\begin{cases} u_{tt} - a^2 u_{xx} - c u_t = 0, \\ u(x, 0) = 0, \\ u_t(x, 0) = 0 \end{cases}$$

⑥ We see ~~that~~ from the first part
of the question that

$$\frac{d}{dt} E[u](t) = 0,$$

But $E[u](0) = \int_0^2 f_0^2 dx = 0$

So $E[u](t) \equiv 0$ for all t .

Namely $U_t \equiv 0, U_x \equiv 0$

and so $U \equiv 0$,

we must have $U_1 = U_2$

and the solution to

$$\begin{cases} U_{tt} - a^2 U_{xx} - c U_t = f(x), \\ U(x, 0) = f(x), \quad U_t(x, 0) = g(x) \end{cases}$$

is unique.

PS 6 Q1:

using the formula (3.6) in Week 6 notes
for inhomogeneous problem with

$$u(x) = \sin x$$

$$f(x) = 0$$

$$g(x) = \sin x$$

and $c = 1$

we get

$$\begin{aligned}
 u(x,t) &= \frac{0+0}{2} + \frac{1}{2} \int_{t-t}^{x+ct} \sin s ds \\
 &\quad + \frac{1}{2} \int_0^t \int_{t+s}^{x+ct+s} \sin r dr ds \\
 &= \frac{\cos(x-ct) - \cos(x+ct)}{2} \\
 &\quad + \frac{1}{2} \int_0^t [\cos(x-t+s) - \cos(x+t-s)] ds \\
 &= \frac{\cos(x-ct) - \cos(x+ct)}{2} + \frac{\sin x - \sin(x+t)}{2}
 \end{aligned}$$

$$- \frac{1}{2} \cdot [-\sin x + \sin(x+t)]$$

$$= \frac{\cos(x-ct) - \cos(x+ct)}{2} + \sin x - \frac{\sin(x-ct) + \sin(x+ct)}{2}$$

PS 6 Q4:

For this question, we can directly apply
the formula in section 3 of Week 5 notes

with $L = \pi$

$$f = b \sin x + 2b \sin(2x)$$

$$g = \sin x$$

thus $a_n = \frac{2}{\pi} \int_0^\pi [b \sin x + 2b \sin(2x)] \sin(nx) dx$

$$b_n = \frac{2}{n\pi c} \int_0^\pi \sin x \sin(nx) dx$$

Using the orthogonality of $\sin nx$ and $\sin mx$
we see all $a_n = 0$ except for $n=1, 2$,

and $a_1 = b$, $a_2 = 2b$

all $b_n = 0$ except for $n=1$

and $b_1 = \frac{1}{c}$

so

$$u(x-t) = b \cdot \sin x \cos(ct) + 2b \sin(2x) \cos(2ct) \\ + \frac{1}{c} \sin x \sin(ct)$$

PS 6 Q 8:

- We differentiate $E(U(t))$ with respect to t and get

$$\begin{aligned}
 & \frac{d}{dt} E(U(t)) \\
 = & \frac{d}{dt} \frac{1}{2} \int_{-\infty}^{\infty} [U_t^2 + \frac{1}{2} U_{xx}^2] dx \\
 = & \frac{1}{2} \int_{-\infty}^{\infty} \frac{d}{dt} [U_t^2 + \frac{1}{2} U_{xx}^2] dx \quad \checkmark \text{chain rule} \\
 = & \frac{1}{2} \int_{-\infty}^{\infty} [2U_t \cdot U_{tt} + 2\frac{1}{2} U_x \cdot U_{xt}] dx \\
 = & \int_{-\infty}^{\infty} U_t \cdot U_{tt} dx + \frac{1}{2} \int_{-\infty}^{\infty} U_x \cdot U_{xt} dx \quad \checkmark \text{integration by parts} \\
 = & \int_{-\infty}^{\infty} U_t \cdot U_{tt} dx + \frac{1}{2} U_x \Big|_{x=-\infty}^{\infty} - \frac{1}{2} \int_{-\infty}^{\infty} U_{xx} \cdot U_t dx \\
 = & \int_{-\infty}^{\infty} U_t \cdot U_{tt} dx - \int U_t \cdot \frac{1}{2} U_{xx} dx \quad \checkmark \text{using compact support.} \\
 = & \int_{-\infty}^{\infty} U_t [U_{tt} - \frac{1}{2} U_{xx}] dx
 \end{aligned}$$

$$\begin{aligned}
 & = \int_{-\infty}^{\infty} U_t \cdot C U_t \quad \checkmark \text{using the PDE} \\
 & \qquad U_{tt} - \frac{1}{2} U_{xx} = C U_t \\
 & = C \int_{-\infty}^{\infty} U_t^2 \\
 & \leq 0 \quad (\text{because } C < 0 \text{ and } \int_{-\infty}^{\infty} U_t^2 \geq 0) \\
 & \text{so the energy is non-increasing.}
 \end{aligned}$$

- For the second part of the question, suppose u_1 and u_2 are 2 solutions to
$$\begin{cases} u_{tt} - \delta^2 u_{xx} - c u_t = f(x), & t < 0 \\ u(x, 0) = g(x), \quad u_t(x, 0) = h(x) \end{cases} \quad (\text{**})$$

then $v = u_1 - u_2$ is a solution to

$$\begin{cases} v_{tt} - \delta^2 v_{xx} - c v_t = 0 \\ v(x, 0) = 0, \quad v_t(x, 0) = 0 \end{cases}$$

Notice the initial condition gives

$$E[v](0) = \frac{1}{2} \int_0^t 0 + \delta^2 \cdot 0 = 0$$

and the first part shows us $E[v](t)$ is non-increasing.

But because $E[v](t) \geq 0$ for all t ,
so we must have $E[v](t) \equiv 0$

Namely $v_x \equiv 0$ and $v_t \equiv 0$.

$$\begin{aligned} v(x, t) &= \int_0^t v_t(x, s) ds + v(x, 0) \\ &= 0 + 0 = 0 \end{aligned}$$

So $v \equiv 0$ and thus $u_1 \equiv u_2$
and the solution to (***) is unique.

Solution to PS 6 Q7:

Show that the energy of the solution $u(x,t)$ to the problem

$$\begin{cases} u_{tt} - c^2 u_{xx} = 0, & x \geq 0, t > 0 \\ u(0,t) = 0 \\ u(x,0) = f(x), \quad u_t(x,0) = g(x) \end{cases} \quad (\star\star)$$

for the wave equation on half line is conserved.

Assuming f, g are compact support, use the conservation of energy to show uniqueness of solutions.

$$\text{The Energy } E[u](t) = \int_0^\infty \left(\frac{1}{2} u_t^2 + \frac{1}{2} c^2 u_x^2 \right) dx$$

has time derivative

$$\frac{d}{dt} E[u](t) = \int_0^\infty \left(\frac{1}{2} \cdot 2u_t u_{tt} + \frac{1}{2} c^2 \cdot 2u_x u_{xt} \right) dx$$

$$\rightarrow = \int_0^\infty u_t \cdot u_{tt} dx + c^2 u_x \cdot u_t \Big|_0^\infty - \int_0^\infty c^2 u_{xx} \cdot u_t dx$$

(Integration by parts)

The boundary condition gives $u_t(0,t) = \frac{d}{dt} u(0,t) = 0$

The condition of compact support gives $u_x \cdot u_t$ is zero at ∞ .

$$\text{so } \frac{d}{dt} E[u](t) = \int_0^\infty u_t \cdot (u_{tt} - c^2 u_{xx}) dx$$

$$\rightarrow = \int_0^\infty u_t \cdot 0 dx$$

by the equation $= 0$

so Energy is conserved.

Next, if u_1 and u_2 are 2 solutions to $(*)$

then $w = u_1 - u_2$ is a solution to

$$\left\{ \begin{array}{l} u_{tt} - c^2 u_{xx} = 0, \quad x \geq 0, t \geq 0. \\ u(0, t) = 0 \\ w(x, 0) = 0, \quad u_t(x, 0) = 0 \end{array} \right.$$

Noticing $E[w](0) = 0$, the preservation of energy implies $E[w](t) = 0$ for all t .

so $w_t \equiv 0$, and thus

$$w(x, t) = \int_0^t w_t(x, t') dt' + w(x, 0) = 0 + 0 \equiv 0$$

Thus $u_1 - u_2 \equiv 0$ and $u_1 \equiv u_2$.

The solution is unique.

