

DYNAMICAL SYSTEMS

MTH744 U/P

SEMESTER A

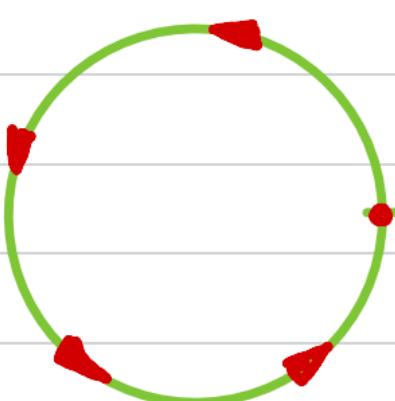
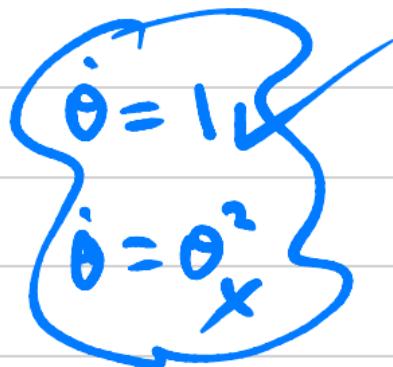
WEEKS

5, 6, 7, 8, 9

2023 – 2024

David Arrowsmith

3. Dynamics on the Circle S^1



$$\dot{\theta} = 1 - \cos \theta$$

D=0 FPs

$$\dot{\theta} = 0$$

$$\text{i.e. } 1 - \cos \theta = 0$$

$$\theta = 0 \quad \text{FP}$$

note $\dot{\theta} > 0, \dot{\theta} \neq 0$

$$\dot{\theta} = f(\theta)$$

needs $f(\theta + 2\pi) = f(\theta)$

Dynamics system on the circle

$$\frac{d\theta}{dt} = \dot{\theta} = f(\theta), \quad \theta \in S^1$$

c.f. $\dot{x} = f(x)$

$$\dot{x} = f(x)$$

$gr(f)$

For $\dot{\theta} = f(\theta)$, we need periodic f e.g.

representing
if $f(x)$ as a
graph $gr(f)$



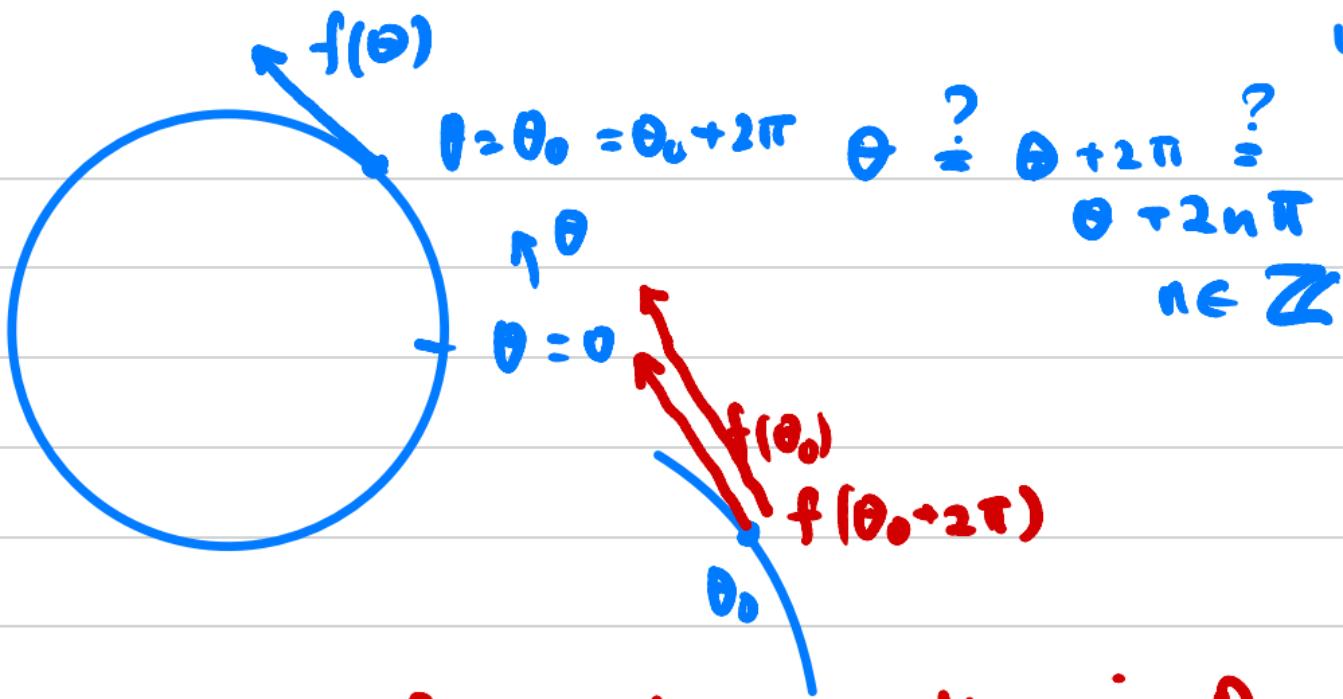
avoids overlaying
vectors horizontally



"sin 30"



W5.2



We need f to be periodic in θ

$$\text{i.e. } f(\theta) = f(\theta + 2k\pi), k \in \mathbb{Z}$$

Fourier series

$$f(\theta) = \sum_{k \in \mathbb{Z}} a_k \cos(kx) + b_k \sin(kx)$$

p18

$$f(\theta) = \dots$$

$$I \subseteq \mathbb{Z}$$

periodicity 2π

Ex 3.1.

 $\dot{\theta} = \omega \rightarrow \text{uniform motion}$
Note $k=0$

$$f(\theta) = a_0 \cdot \checkmark$$

$$\frac{d\theta}{dt} = \omega \quad \theta = \omega t + \theta_0$$

$$, \quad t=0 \\ \theta=\theta_0$$

constant speed
not constant vel.

cf. $\dot{x} = w$
on \mathbb{R}

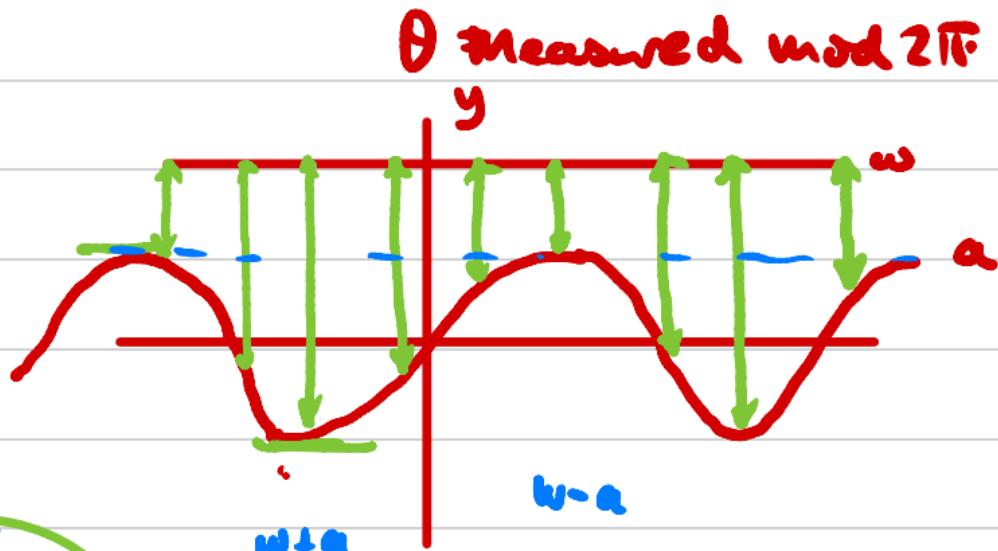
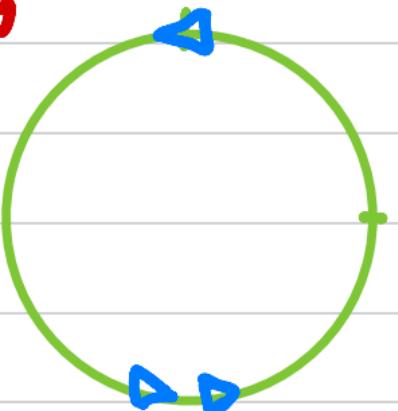
const speed
& velocity

$$\dot{\theta} = \omega - a \sin \theta.$$

$$\omega, a > 0$$

$$\omega > a > 0$$

$$\omega - a \sin \theta$$



at $\theta = \pi/2$ or $a \rightarrow \omega$

speed slows to zero

$$\theta = -\pi/2$$

Speed increases to 2ω

What happens as $a \rightarrow \omega^-$

Slows down $\frac{T\pi}{2}$
Speeds up at $\frac{T\pi}{2}$

$$a = \omega$$

$$\dot{\theta} = \omega(1 - \sin \theta)$$

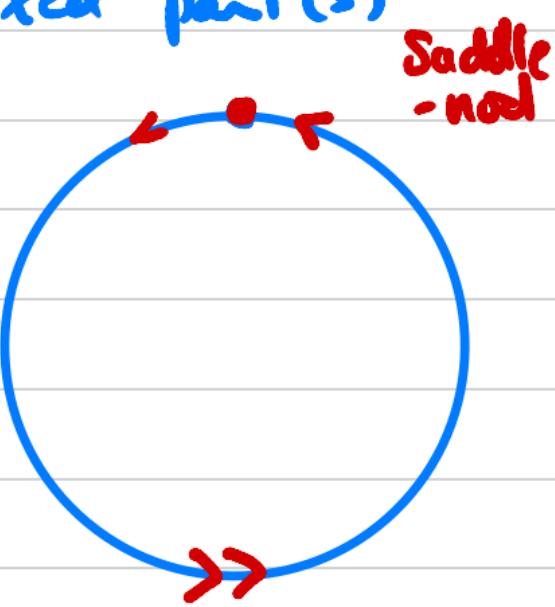
$$\sin \theta = 1, \quad \theta = \frac{\pi}{2}$$

\exists fixed point(s)

Saddle
-node

$$\int dt = \int_{\theta_0}^{\theta} \frac{d\theta}{\omega - a \sin \theta}$$

P19



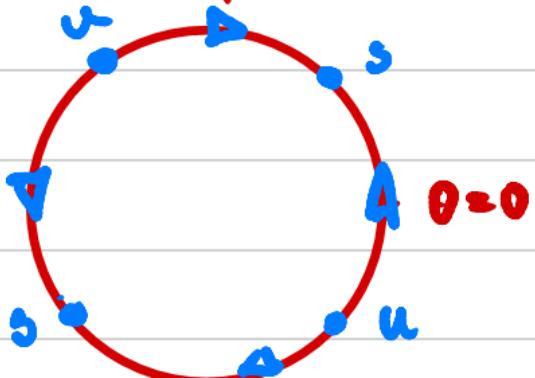
As $a \rightarrow w^-$, then T for a complete circle to be covered by an orbit (solution arc) tends to ∞ .

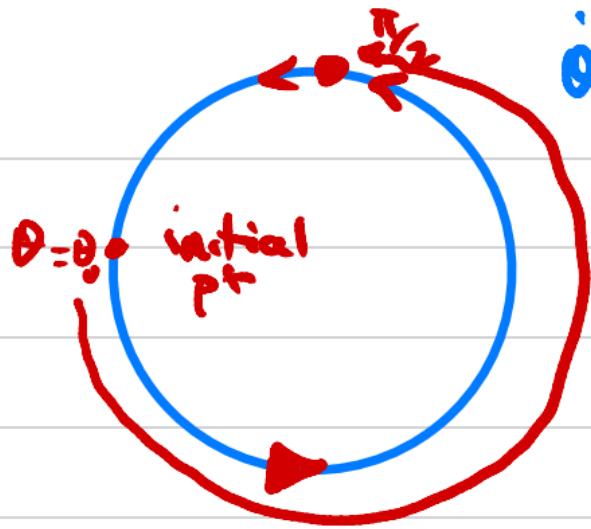
As $a \rightarrow w^-$, a speed bottle-neck is created around $\Theta = \pi/2$, and when we take the limit of $w=a$, a diked part and a complete circuit no longer happens

p19 $T = \frac{2\pi}{\sqrt{\omega^2 - a^2}}$ $\Rightarrow a \rightarrow w' T \rightarrow \infty$

Ex $\dot{\theta} = \cos 2\theta$ FPs $\cos 2\theta = 0$
 $2\theta = \pi/2, -\pi/2, \frac{5\pi}{2}, +\frac{3\pi}{2}$

$$\theta = \frac{\pi}{4}, -\frac{\pi}{4}, \frac{5\pi}{4}, \frac{3\pi}{4}$$





$$\dot{\theta} = \omega(1 - \sin \theta)$$

what set of points
on S^1 are attracted to

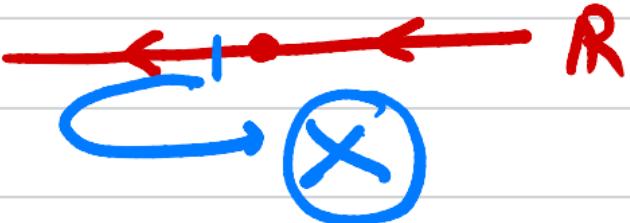
$$\theta = \pi r_2$$

$$\{ \theta_0 \mid \theta(t) \text{ with } \theta(0) = \theta_0 \\ \rightarrow \pi r_2 \text{ as } t \xrightarrow{\nearrow} \infty \}$$

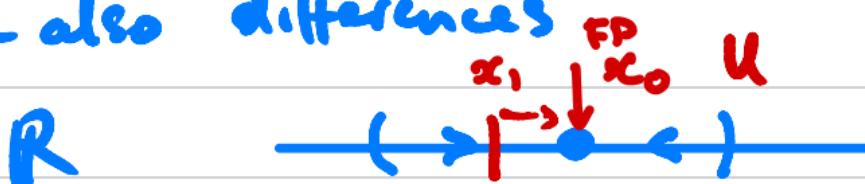
Basin of attraction of the fixed pt $\theta = \pi r_2$

$$B\left(\frac{\pi}{2}\right) = S^1$$

cf.



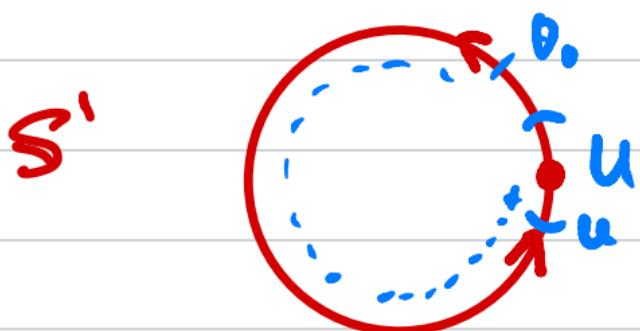
These examples show similarities with real life dynamical systems, but there are also differences



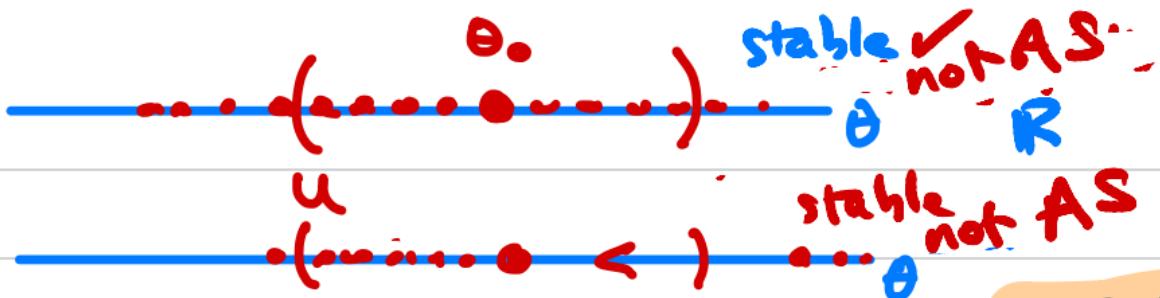
Start in U , nbd of x^*
we stay in U
 $\rightarrow x_0$

x_0 is a stable fixed pt. Let $x_1 \in U$

then $x(t) \rightarrow x_0$, where $x(0) = x_1$. Also P.S. ($x(t) \rightarrow x_0$)
 $\forall x_1 \in U$



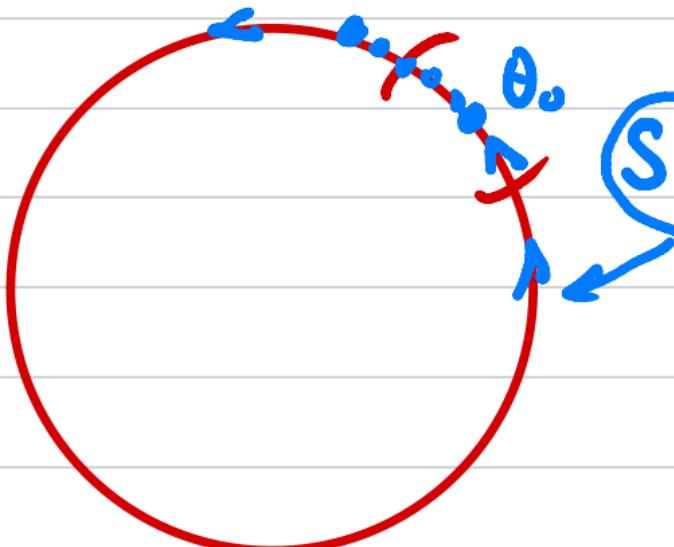
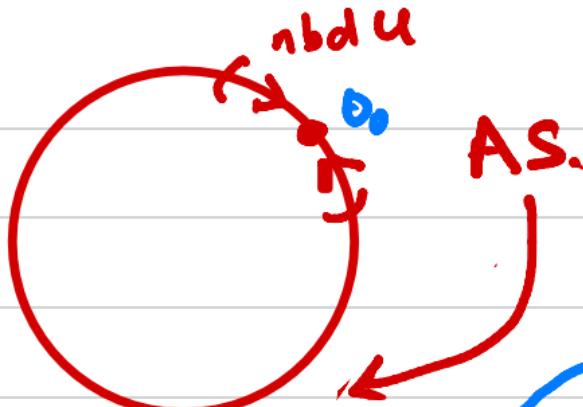
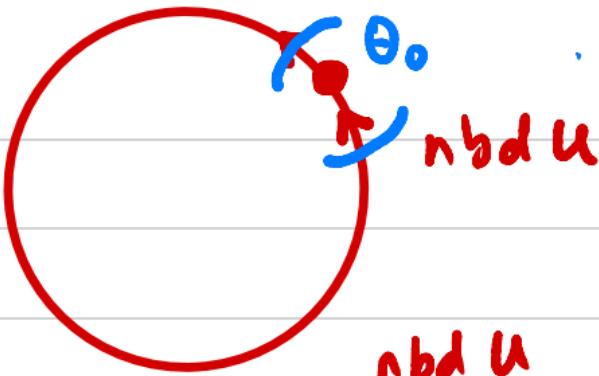
start in U
sometimes we
leave and
then RE-ENTER!
 $\rightarrow D_0$



W5.10

- stability on S^1 at fp $\theta = \theta_0$.
 $\exists U$ s.t. if $\theta_1 \in U$, then orbit $\Theta(t)$
 with $\Theta(0) = \theta_1$ remains in U as $t \rightarrow \infty$

asymptotic stability if $\theta_1 \in U$, $\Theta(t) \xrightarrow{*} \theta_0$
 as $t \rightarrow \infty$.

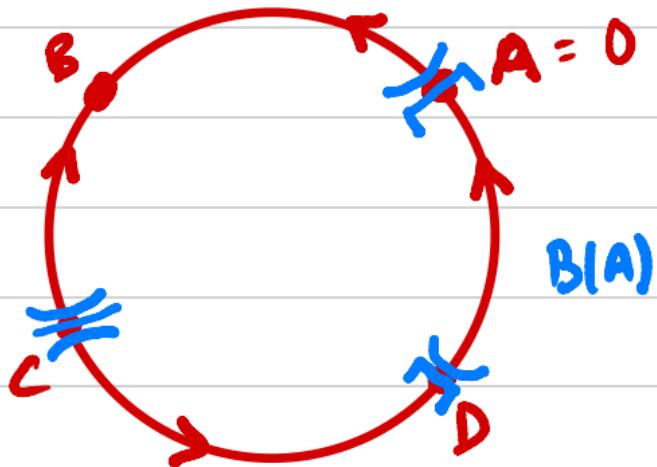


S=stable

AS=asymptotically stable

$\theta = \theta_0$ is a fixed pt in the circle phase portraits

Example of Basins of attraction



$$\bigcup_{x \in \{A, B, C, D\}} B(x) = S^1$$

Basins of attraction partition the state space S . $\dot{x} = x, B(0) = \{0\}, B(\infty) = \{(\infty, 0)\}$
 Note for \mathbb{R} , we need to accommodate infinity

A few calculations on time to rotate a full circle

W5.13

$$\dot{\theta} = \omega - a \sin \theta \Rightarrow \int \frac{d\theta}{\omega - a \sin \theta} = \int dt$$

$$u = \tan\left(\frac{\theta}{2}\right) \Rightarrow du = \sec^2 \theta d\theta \Rightarrow du = \frac{(1+u^2)d\theta}{2}$$

$$\sin \theta = 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2} = 2 \tan \frac{\theta}{2} \cdot \cos^2 \frac{\theta}{2} = \frac{2u}{1+u^2}$$

$$\int \frac{d\theta}{\omega - a \sin(\theta)} = \int \frac{2du}{(1+u^2)(\omega - a \frac{2u}{1+u^2})} = \int \frac{2du}{\omega(1+u^2) - 2au}$$

$$= \frac{1}{\omega} \int \frac{2du}{u^2 - 2\frac{a}{\omega}u + 1} = \frac{1}{\omega} \int \frac{2du}{\left(u - \frac{a}{\omega}\right)^2 + \left(1 - \frac{a^2}{\omega^2}\right)} =$$

$$\frac{2}{\omega} \frac{1}{\sqrt{1-\frac{a^2}{\omega^2}}} \tan^{-1} \left(\frac{u - \frac{a}{\omega}}{\sqrt{1-\frac{a^2}{\omega^2}}} \right)$$

$$\begin{cases} \theta = \pi, u = \tan \frac{\pi}{2} = \infty \\ \theta = -\frac{\pi}{2}, u = -\tan \frac{\pi}{2} = -\infty \end{cases}$$

$$\frac{2}{\sqrt{\omega^2 - a^2}} \left[\tan^{-1} \left(\frac{u - \frac{a}{\omega}}{\sqrt{1 - \frac{a^2}{\omega^2}}} \right) \right]_{-\infty}^{\infty} = \frac{\pi}{\sqrt{\omega^2 - a^2}}$$

W5.44

See Latex notes 'if' interested.

Dynamical System on \mathbb{R}^2 (introduction)

Linear Systems on \mathbb{R}^2

Ch 4

→ Linear Systems on \mathbb{R}

$$\dot{x} = \underline{ax}$$

Linearity cond

$$f(a)x + b x' = af(x) + bf(x')$$

$$a > 0$$



$$a = 0$$



$$a < 0$$



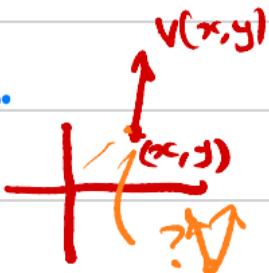
line of
fixed pts

That's it!!

Chapter 4

Linear system in \mathbb{R}^2

$$\begin{cases} \dot{x} = ax + by \\ \dot{y} = cx + dy \end{cases} \quad a, b, c, d \in \mathbb{R}$$



$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

vector field at the point (x, y) is
 $v(x, y) = (ax + by, cx + dy)$

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} a & 0 \\ 0 & d \end{bmatrix}$$

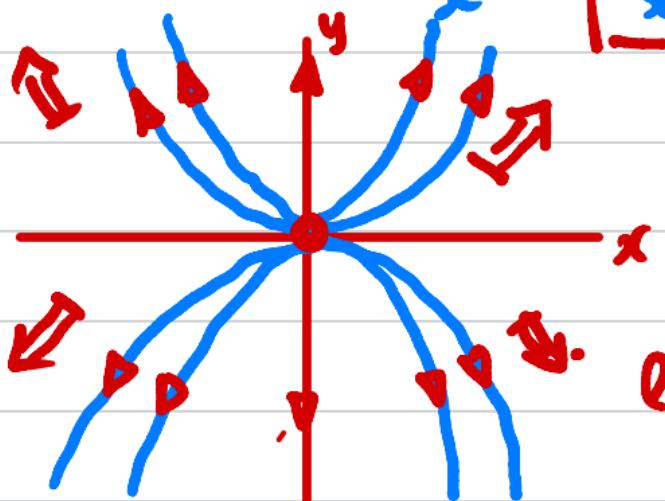
- consider this simple case (diagonal) matrix

$$\begin{aligned}\dot{x} &= ax \\ \dot{y} &= dy\end{aligned}\quad (b=c=0)$$

$$\begin{aligned}\dot{x} &= x \\ \dot{y} &= 2y\end{aligned}$$

FP $\begin{cases} \dot{x}=0 \\ \dot{y}=0 \end{cases} \quad \begin{cases} x=0 \\ y=0 \end{cases}$

elim dt $\frac{\dot{y}}{\dot{x}} = \boxed{\frac{2y}{x} = \frac{dy}{dx}}$



$$\int 2 \frac{dx}{x} = \int dy$$

$$2 \ln x + c = -\ln y$$

$$\begin{aligned}2 \ln x + 2 \ln(y) + c &= -\ln(A) \\ \ln A + 2 \ln(x) + 2 \ln(y) &= 0 \\ y &= Ax^2\end{aligned}$$

"unstable node"

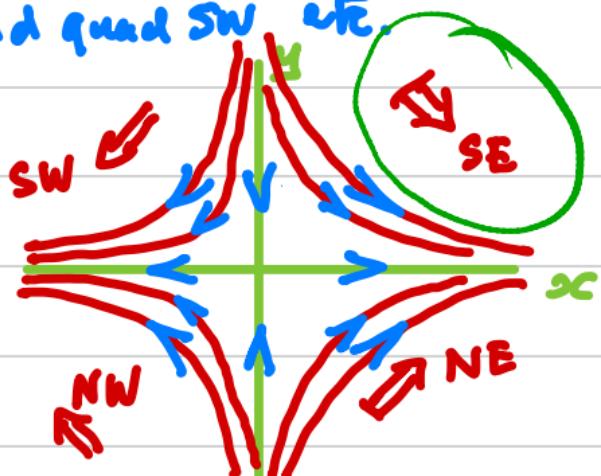
wb.4

$$\dot{x} = x$$

$$\dot{y} = -y$$

1st quad $x, y > 0$,
 $\therefore \dot{x} > 0$ "SE"
 $\dot{y} < 0$

2nd quad SW etc.



$$\frac{dy}{dx} = -\frac{y}{x}$$

$$\frac{dy}{y} + \frac{dx}{x} = 0$$

$$\ln(y) + \ln(x) = C$$

$yx = C'$, hyperbolae

phase portrait of a saddle.

LECTURES WEEK 6

W6.5

Linear Systems

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix}$$

$$\dot{\underline{z}} = A\underline{z},$$

$$A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

$$\underline{w} = \begin{bmatrix} u \\ v \end{bmatrix}$$

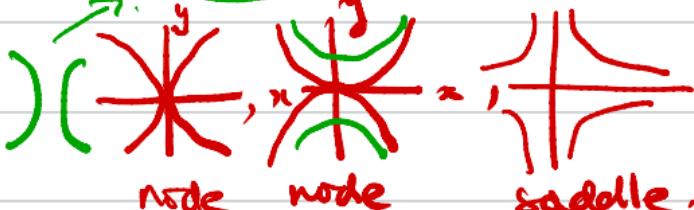
Simple examples of $A = \text{diagonal matrix}$ last time

$\dot{x} = ax, \dot{y} = y$ for various values of a

Simple to work out $\frac{dy}{dx} = \frac{dy}{dx} = \frac{ay}{y} = a \Rightarrow x = Cy^a$

$a > 1, a < 1$

, $a < 0$ etc.



$$\boxed{\dot{x} = y, \dot{y} = -\omega^2 x \text{ for } y \text{-axis}}$$

$$\boxed{H = y^2 + \omega^2 x^2 \text{ for } x \text{-axis}}$$

How does H vary with time?

Consider:

$$\dot{x} = y, \quad \dot{y} = -\omega^2 x, \quad ,$$

$$H = \omega^2 x^2 + y^2$$

$$H(t) = \omega^2 x(t)^2 + y(t)^2$$

$$\frac{dH(t)}{dt} = 2\omega^2 x \dot{x} + 2y \dot{y} =$$

$$= 2\omega^2 x y - 2y \omega^2 x \equiv 0$$

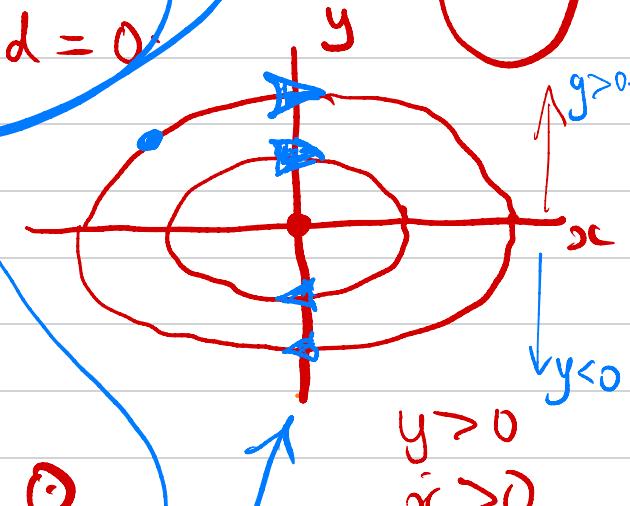
$$H(t) = \text{constant} \cdot \mathcal{E}: \lambda^2 + \omega^2 = 0 \quad \text{periodic orbits.}$$

$$\lambda = \pm i\omega$$

$$a=0 \quad b=1$$

$$c=-\omega^2 \quad d=0$$

$dt?$



Eigenvalues of $A = \{\lambda_1, \lambda_2\}$.

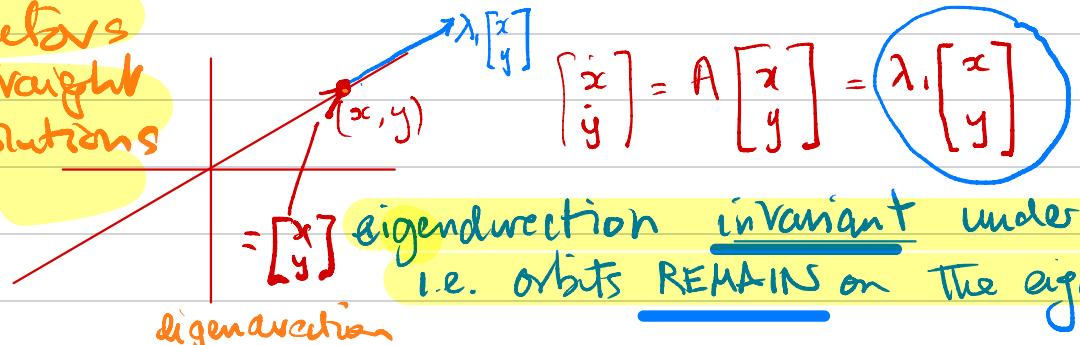
Eigenvectors of A $v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, v_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$

$$A = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \quad \lambda_1 \neq \lambda_2$$

Note in linear systems {Solution curves
Trajectories lie along the
orbits} eigendirection (in general)

Suppose $A \begin{bmatrix} x \\ y \end{bmatrix} = \lambda_1 \begin{bmatrix} x \\ y \end{bmatrix}$, v_1 = eigenvector of λ_1 ,

eigenvectors
give straight
line solutions
orbits.



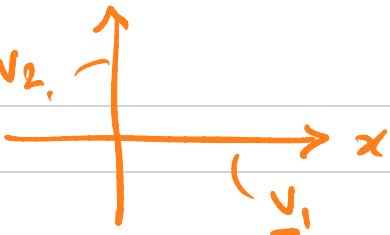
eigendirection invariant under the flow
i.e. orbits REMAINS on the eigendirection.

$$\begin{aligned} v_1 &= \begin{bmatrix} u \\ v \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ A \begin{bmatrix} u \\ v \end{bmatrix} &= \lambda_1 \begin{bmatrix} u \\ v \end{bmatrix} \end{aligned}$$

$$\begin{aligned} \lambda_1 u &= \lambda_1 u \\ \lambda_2 v &= \lambda_1 v \end{aligned}$$

$$\lambda_1 \neq \lambda_2$$

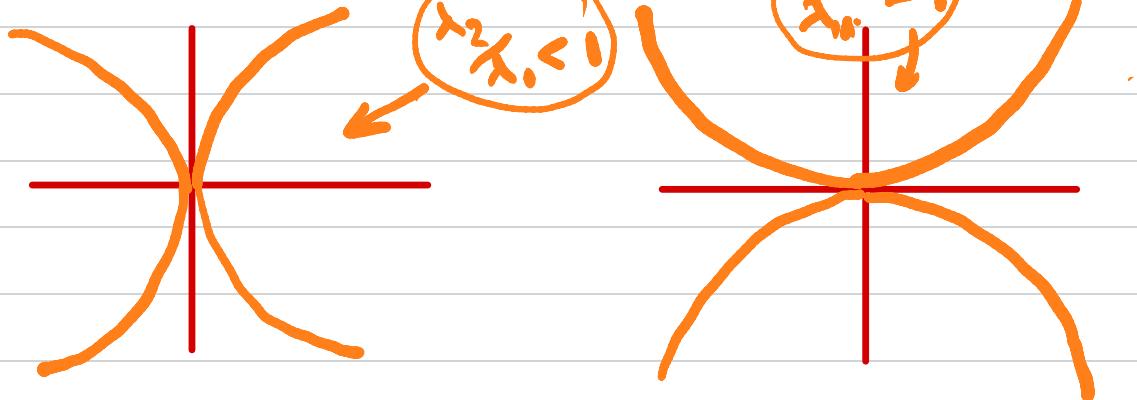
$$A = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}$$



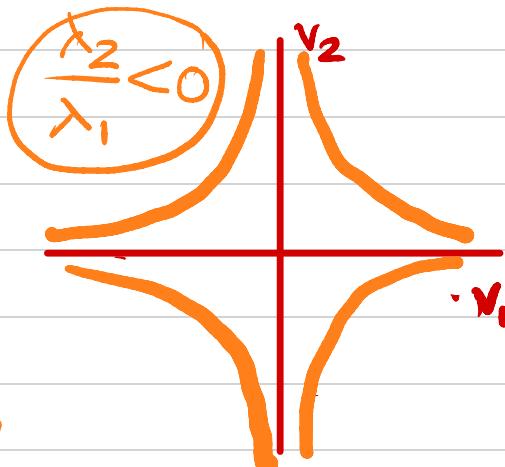
$$\frac{dx}{\lambda_1 x} = \frac{dy}{\lambda_2 y}$$

~~$= dt$~~

$$y = C x^{\frac{\lambda_2}{\lambda_1}}$$



$$\frac{\lambda_2}{\lambda_1} > 1$$



Jordan forms

$$\dot{\underline{z}} = A \underline{z}$$

$$, \underline{z} = \begin{bmatrix} x \\ y \end{bmatrix}$$

$$\underline{w} = \begin{bmatrix} u \\ v \end{bmatrix} \text{ - new coordinates}$$

Let $\underline{z} = P \underline{w}$ P is non-singular $\underline{P}^{-1} \exists$

$$\dot{\underline{z}} = A \underline{z}, \quad \dot{\underline{z}} = P \dot{\underline{w}}$$

$$\underline{z} = \begin{bmatrix} pu + qv \\ ru + sv \end{bmatrix}$$

$$(P \dot{\underline{w}}) = A \underline{z} = (A P \underline{w})$$

$$P = \begin{bmatrix} p & q \\ r & s \end{bmatrix}$$

$$\dot{\underline{w}} = P^{-1} A P \underline{w}$$

Can I choose P such
that $P^{-1} A P$ is "simpler" than A .

$$= J \underline{w}$$

constants

A , calculate eigenvalues λ_1, λ_2

$$J = P^{-1} A P \text{ similar}$$

λ_1, λ_2 real and $\lambda_1 \neq \lambda_2$

$$\exists P, P^{-1} A P = \begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ 0 & & \lambda_2 \end{bmatrix} \quad \checkmark$$

$$E(J) \stackrel{\text{to } A}{=} E(A)$$

λ_1, λ_2 real and $\lambda_1 = \lambda_2$ but A is not diagonal.

$(=\lambda)$

$$\exists P, P^{-1} A P = \begin{bmatrix} \lambda & & \\ & \ddots & \\ 0 & & \lambda \end{bmatrix}$$

λ_1, λ_2 real $\lambda_1 = \lambda_2$ A is diagonal.

$$A = \begin{bmatrix} \lambda & 0 \\ 0 & \lambda \end{bmatrix} \quad (= \lambda)$$

$$\exists P = I \text{ s.t. } P^{-1} A P = J \not\models A$$

λ_1, λ_2 complex: $\lambda_1 = \alpha + i\beta$ $\lambda_2 = \alpha - i\beta$

$(A \text{ is real})$

w6.11

$$\exists P \text{ (real)} \quad \text{s.t.} \quad P^{-1}AP = \begin{bmatrix} \alpha & -\beta \\ \beta & \alpha \end{bmatrix}$$

$$[\exists Q \text{ s.t. } Q^{-1}AQ = \begin{bmatrix} \alpha+i\beta & 0 \\ 0 & \alpha-i\beta \end{bmatrix}] \quad \beta \neq 0$$

↑
Real → Complex

$$i = (\alpha + i\beta) u, \quad j = (\zeta - i\beta) v$$

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} \alpha - \beta \\ \beta \end{bmatrix} \cdot \begin{bmatrix} x \\ y \end{bmatrix}$$

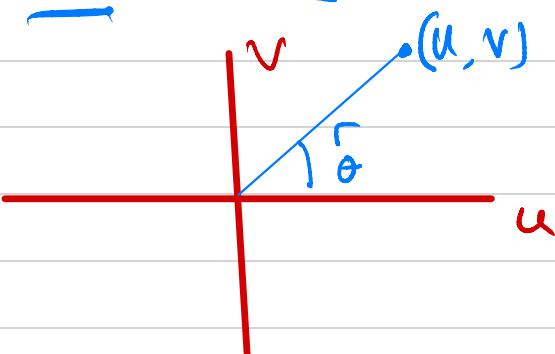
$$\underline{w} = J \underline{u}$$

Polar coordinates

$$x = r \cos \theta$$

$$y = r \sin \Theta$$

A - level.



$$r^2 = x^2 + y^2, \quad rr' = x\dot{x} + y\dot{y}$$

P27

$$rr' = uu + vv$$

$$= u(\alpha u - \beta v) + v(\beta u + \alpha v)$$

$$= \alpha(u^2 + v^2)$$

$$rr' = \alpha r^2 \Rightarrow \boxed{\dot{r} = \alpha r}$$

$\alpha > 0$, r incr with time
 $\alpha < 0$, r decr with time

$$rr' = \alpha r^2$$

$$\dot{r} = \alpha r$$

$$\dot{\theta} = \beta$$

$$r^2\dot{\theta} = xy - y\dot{x}$$

"uv coords"

$$r^2\dot{\theta} = uv - vi\dot{u} = \beta(u^2 + v^2) = \rho r^2$$

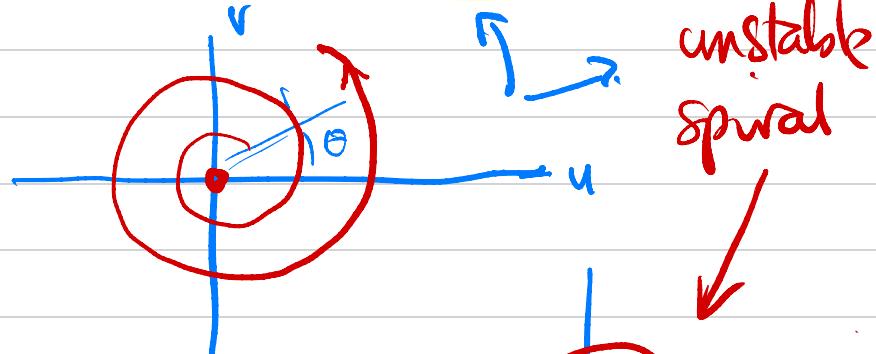
$$\dot{\theta} = \beta$$

Polar coords

General cases of complex eigenvalues

(i) $\alpha > 0$, $\beta > 0$

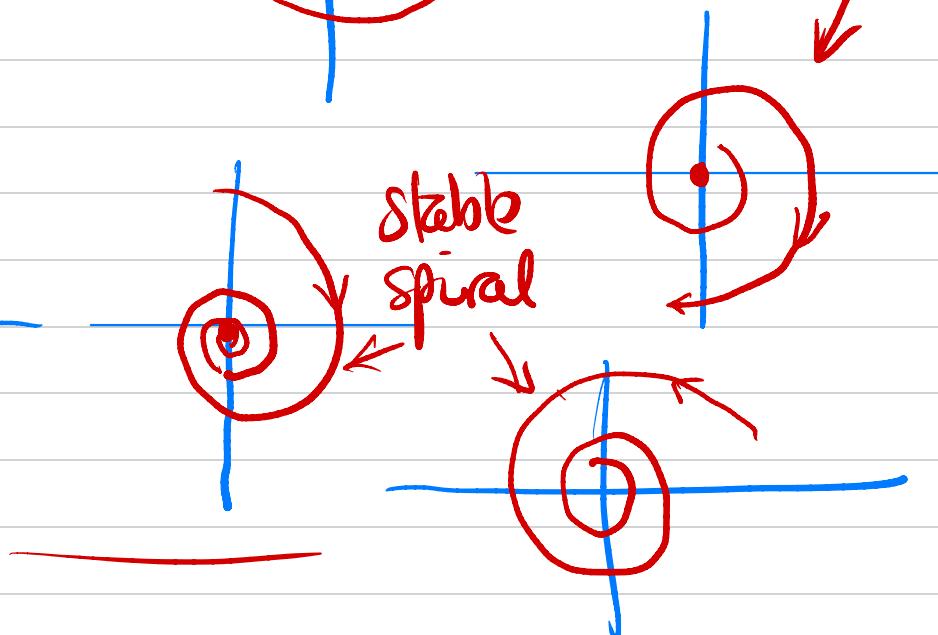
$$r = \alpha r, \theta = \beta$$



(ii) $\alpha > 0$, $\beta < 0$

(iii) $\alpha < 0$, $\beta < 0$

(iv) $\alpha < 0$, $\beta > 0$



Special Cases of complex eigenvalues

(v) $\alpha = 0$

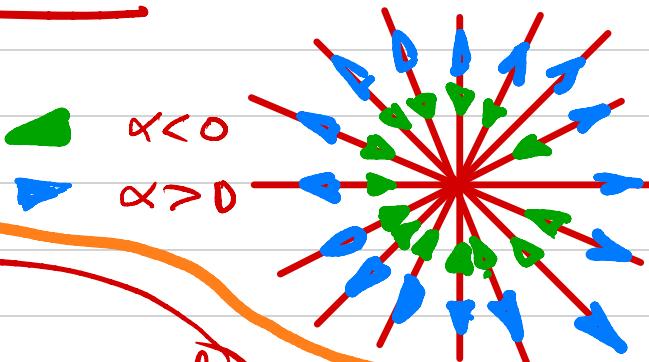
$\beta \neq 0$

\rightarrow Circles

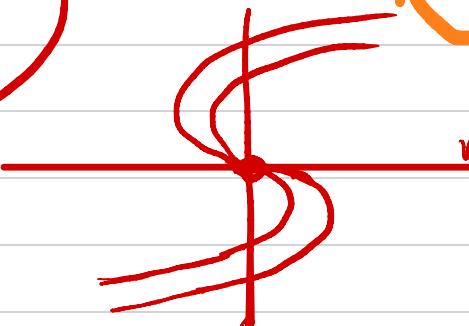
(vi) $\alpha \neq 0$

$\beta = 0$

\rightarrow radial lines



$\lambda_1 = \lambda_2$ non-diagonal
improper node
see Thursday tutorial



earlier example.

$$\dot{x} = y$$

$$\dot{y} = -\omega^2 x$$

$$\beta = \omega.$$

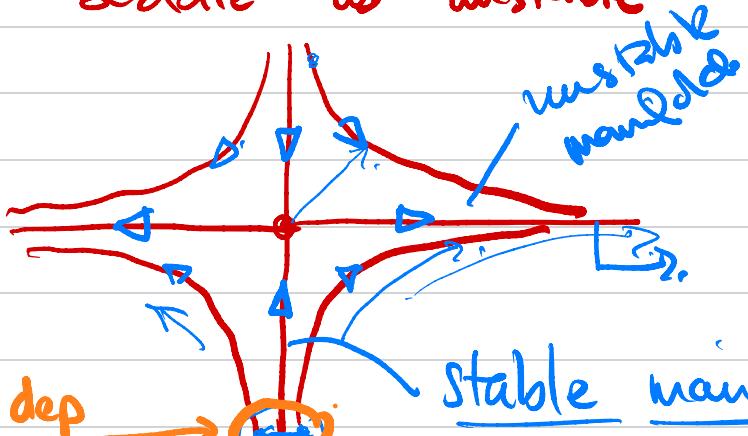
$$\alpha = 0.$$

$\alpha > 0$
unstable
star node

$\alpha < 0$
stable star node

v_1
 v_{-1}
just one eigendirection
because of repeated
eigenvalue

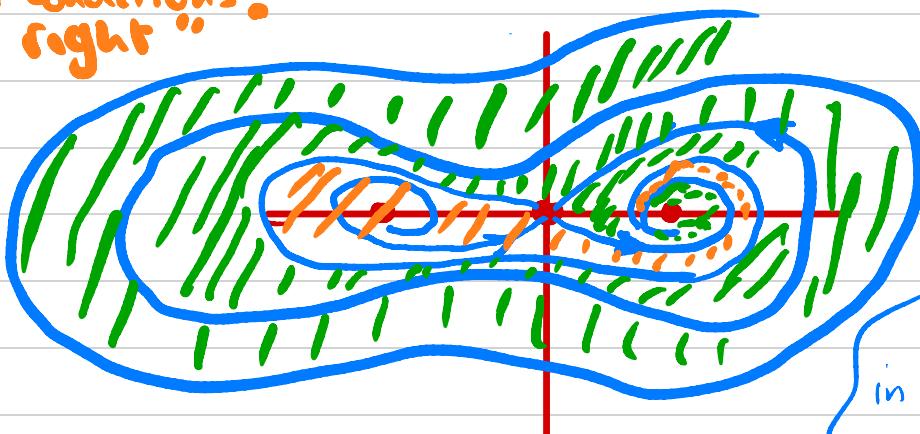
Note Saddle is unstable



$$B(Q) = y\text{-axis}$$

stable manifold of the fixed pt

sensitive dep
on initial conditions!
"left or right"



Nonlinear
System - role
of saddle manifolds
is very important
in dividing basins
of attraction

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix}, \quad \underline{w} = \begin{bmatrix} u \\ v \end{bmatrix} \quad \dot{\underline{z}} = A\underline{z} \rightarrow \dot{\underline{w}} = P^{-1}AP\underline{w}$$

Fig 15 p 30

$$\underline{z} = P\underline{w}$$

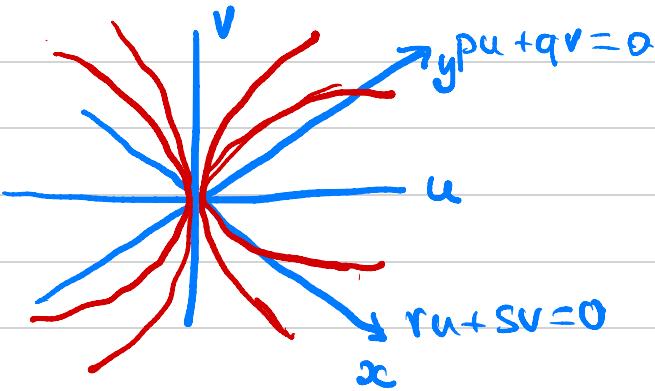
$$\left. \begin{array}{l} x = pu + qv \\ y = ru + sv \end{array} \right\}$$

$$y\text{-axis} \quad x=0$$

$$x\text{-axis} \quad y=0$$

$$x=0 \quad pu + qv = 0$$

$$y=0 \quad ru + sv = 0$$

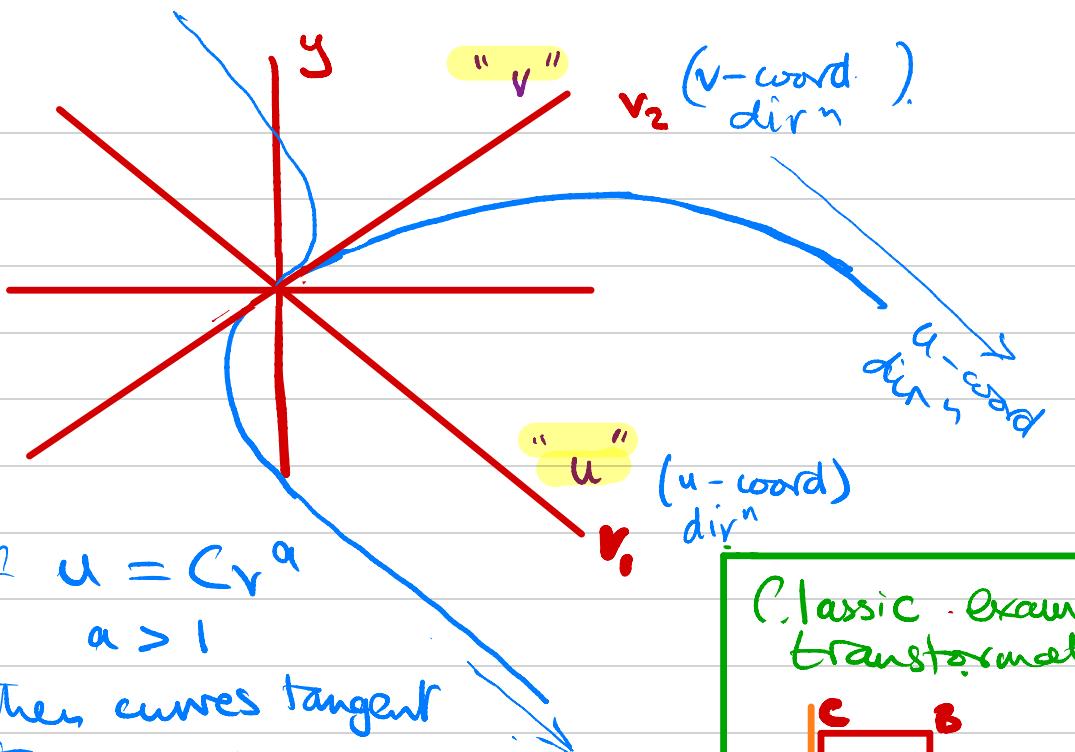


$$\dot{\underline{w}} = \begin{bmatrix} \dot{u} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} 2u \\ v \end{bmatrix}$$

Suppose

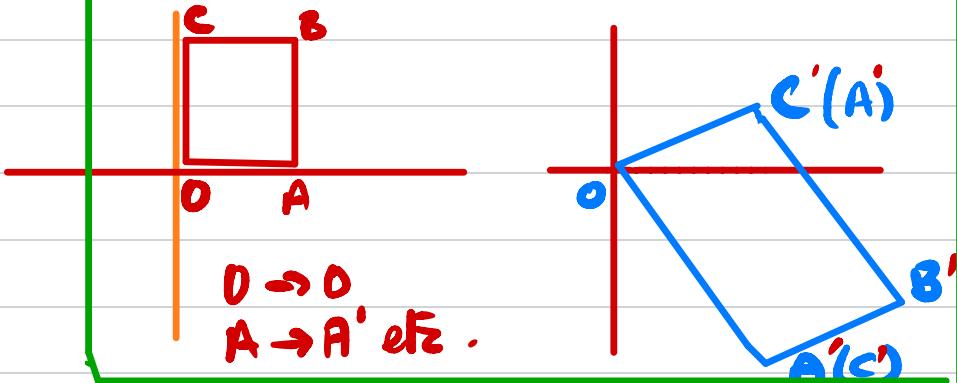
$$\frac{dv}{v} = \frac{du}{2u}$$

$$u = C v^2$$



Σ turn in the u -dirn
 at " ∞ "

Classic example of the linear transformation of a square.



Revision of
week 6.

$$\underline{\dot{z}} = \underline{A} \underline{z}$$

$$\underline{\dot{w}} = \underline{J} \underline{w}$$

Wk8, 1

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix} \approx (x, y)$$

$$\underline{z} = \underline{P} \underline{w}$$

$$\underline{w} = \begin{bmatrix} u \\ v \end{bmatrix}$$

$$\underline{J} = \underline{P}^{-1} \underline{A} \underline{P}$$

$$\approx (u, v)$$

$$\underline{J} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \quad \lambda_1, \lambda_2 \text{ real}$$

$$\begin{array}{l} \lambda_1, \lambda_2 > 0 \\ \lambda_1, \lambda_2 < 0 \end{array}$$

$$\frac{du}{dt} = \lambda_1 u, \quad \frac{dv}{dt} = \lambda_2 v$$

$$\frac{\frac{du}{dt}}{\lambda_1 u} = \frac{\frac{dv}{dt}}{\lambda_2 v}$$

$$\Rightarrow \ln v = \frac{\lambda_2}{\lambda_1} \ln(u) + C$$

$$v = C u^{\frac{\lambda_2}{\lambda_1}}$$

$$v = Cu^{1/2}$$

$$\Rightarrow u = C^2 v^2$$

W8.2a

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = A \begin{bmatrix} x \\ y \end{bmatrix}$$

A has λ_1, λ_2 eigenvalues

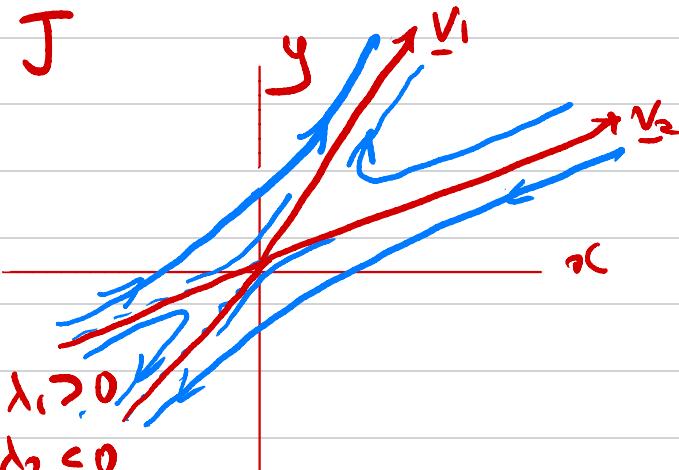
$$\lambda_1, \lambda_2 < 0$$

$$v = Cu^2$$

$v = Cu^2$

$\dot{u} = \lambda_1 u$

$v = \lambda_2 v$



Recap - linear systems and coordinates.

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix}, \underline{w} = \begin{bmatrix} u \\ v \end{bmatrix}, z = \underline{P}\underline{w}$$

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix} = \underline{A} \begin{bmatrix} x \\ y \end{bmatrix} = \underline{A}\underline{z} = \underline{A}\underline{P}\underline{w}$$

$$\underline{P}\underline{w} \quad \therefore \quad \underline{P}\underline{w} = \underline{A}\underline{P}\underline{w}, \underline{w} = \underline{P}^{-1}\underline{A}\underline{P}\underline{w}, \underline{w} = \underline{J}\underline{w}$$

- we choose \underline{P} to give a Jordan matrix \underline{J}
- remember $\mathcal{E}(\underline{J}) = \mathcal{E}(\underline{A})$ because \underline{A} & \underline{J} are "SIMILAR" matrices.

$$\boxed{\det(\lambda\underline{I} - \underline{J})} = \det(\lambda\underline{I} - \underline{P}^{-1}\underline{A}\underline{P}) = \det(\underline{P}^{-1}(\lambda\underline{I} - \underline{A})\underline{P}) = \cancel{\det(\underline{P}^{-1})} \det(\lambda\underline{I} - \underline{A}) \cancel{\det(\underline{P})}$$

$$= \boxed{\det(\lambda\underline{I} - \underline{A})} \quad \boxed{\det(\underline{P}^{-1}) \cdot \det(\underline{P}) = \det(\underline{P}\underline{P}^{-1}) = 1}$$

Change of coordinates

$$\begin{bmatrix} x \\ y \end{bmatrix} \leftrightarrow \begin{bmatrix} u \\ v \end{bmatrix}$$

\underline{x} \underline{w}

Consider the case of $J = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}$ λ_1, λ_2 real
and $\lambda_1 \neq \lambda_2$

Eigenvectors: \underline{v}_1 for λ_1 , \underline{v}_2 for λ_2 $A\underline{v}_1 = \lambda_1 \underline{v}_1$

$$A\underline{v}_2 = \lambda_2 \underline{v}_2$$

Choose $P = \begin{bmatrix} \underline{v}_1 & \underline{v}_2 \end{bmatrix}$

$$\frac{P}{P} \cdot J$$

$$\underline{AP} = A \begin{bmatrix} \underline{v}_1 & \underline{v}_2 \end{bmatrix} = [A\underline{v}_1 \mid A\underline{v}_2] = [\lambda_1 \underline{v}_1 \mid \lambda_2 \underline{v}_2] = \begin{bmatrix} \underline{v}_1 & \underline{v}_2 \end{bmatrix} \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}$$

$$P^{-1}AP = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} = J \quad \underline{\underline{=}}$$

$$\underline{z} = \underline{P} \underline{w} \quad \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} \underline{v}_1 & \underline{v}_2 \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}$$

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} u \underline{v}_1 + v \underline{v}_2 \\ 1 \end{bmatrix}$$

↓ coord. words

$$x \underline{e}_1 + y \underline{e}_2$$

$$\underline{e}_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \underline{e}_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$x \begin{bmatrix} 1 \\ 0 \end{bmatrix} + y \begin{bmatrix} 0 \\ 1 \end{bmatrix} = u \begin{bmatrix} \underline{v}_1 \\ \underline{v}_2 \end{bmatrix} + v \begin{bmatrix} \underline{v}_2 \\ \underline{v}_1 \end{bmatrix}$$

Significance of change of coordinates -

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix} = P \begin{bmatrix} u \\ v \end{bmatrix}$$

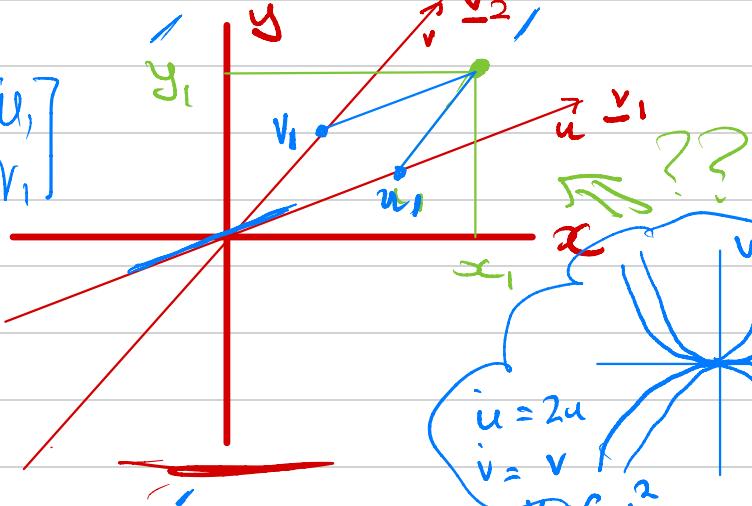
Chosen $P = [v_1; v_2]$ $\begin{bmatrix} x \\ y \end{bmatrix} = [v_1; v_2] \begin{bmatrix} u \\ v \end{bmatrix} = u v_1 + v v_2$

$$\begin{bmatrix} x \\ y \end{bmatrix} = x \begin{bmatrix} 1 \\ 0 \end{bmatrix} + y \begin{bmatrix} 0 \\ 1 \end{bmatrix} = u v_1 + v v_2$$

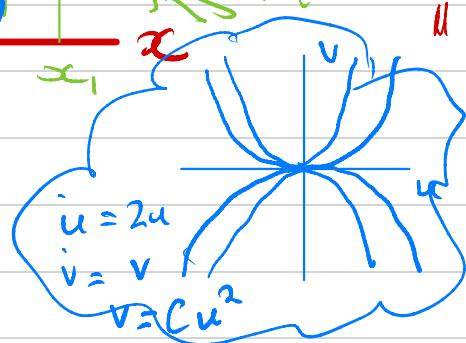
✓✓

u, v coordinates

$$\begin{bmatrix} x_1 \\ y_1 \end{bmatrix} = \begin{bmatrix} v_1 & v_2 \\ u_1 & u_2 \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}$$



$$\begin{bmatrix} u \\ v \end{bmatrix} = \begin{bmatrix} v_1 & 0 \\ 0 & v_2 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$



$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}$$

WK8.5(a)

Note effects of a linear transformation:



WK8.6

Example 4.4 page 32

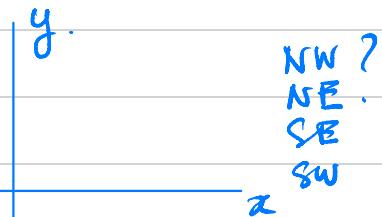
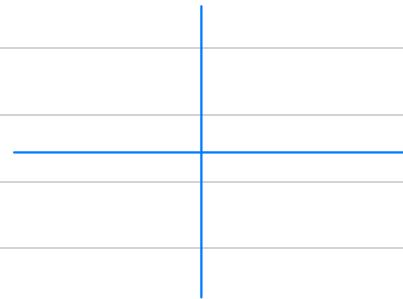
Node

$$A = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix} \quad \checkmark, \quad A = \begin{bmatrix} 4 & -1 \\ 1 & 2 \end{bmatrix}$$

"Jordan canonical form"

$$\lambda_1 = \lambda_2 = 2, \quad \underline{v}_1 = \underline{v}_2 = (1, 0)^T$$

Improper node

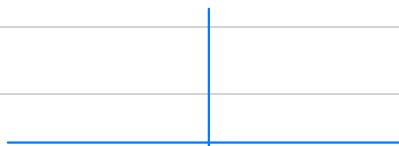


NW?
NE.
SE
SW

$$\begin{aligned} A \underline{v}_1 &= 2 \underline{v}_1 & \underline{v}_1 &= \begin{bmatrix} u \\ v \end{bmatrix} \\ 2u + v &= 2u & \\ 2v &= 2v & \\ v &= 0 & \\ \underline{v}_1 &= \begin{bmatrix} 1 \\ 0 \end{bmatrix} \end{aligned}$$

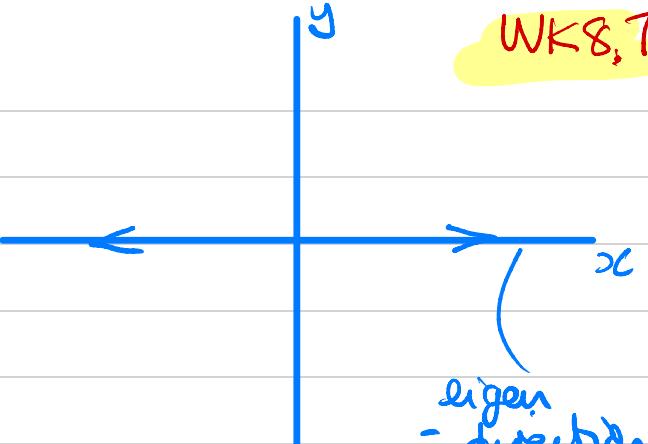
$$\underline{x} = A \underline{z}$$

$$\underline{z} = \begin{bmatrix} x \\ y \end{bmatrix}.$$



$$A = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix}$$

WK8, 7



$$\lambda = \frac{\text{Tr}(A) \pm \sqrt{\text{Tr}(A)^2 - 4 \det(A)}}{2}$$

$$A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

$$\det \begin{bmatrix} \lambda - a & -b \\ -c & \lambda - d \end{bmatrix} = 0$$

"invariant under the flow"
or "if on the eigendirection,
you stay on it."

$$(\lambda - a)(\lambda - d) - bc = 0$$

$$\frac{\lambda^2 - (\text{Tr}(A))\lambda + \det(A)}{\det(A)} = 0$$

$$\lambda = \frac{\text{Tr}(A) \pm \sqrt{\text{Tr}(A)^2 - 4 \det(A)}}{2}$$

2.

$$\text{Tr}(A)^2 - 4 \det(A) > 0$$

$$\text{Tr}(A)^2 - 4 \det(A) = 0$$

$$\text{Tr}(A)^2 - 4 \det(A) < 0$$

nodes/saddles

2 real eigenvalues
 $\lambda_1 \neq \lambda_2$

?? 2 real eigenvalues
 $\lambda_1 = \lambda_2$

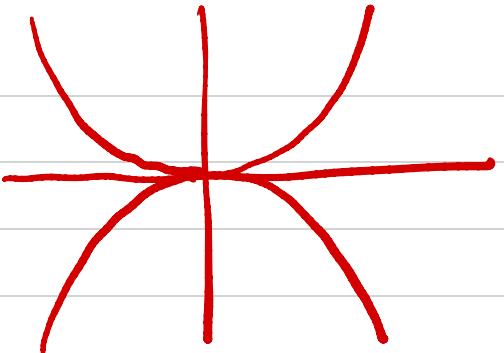
2 complex eigenvalues
 spirals

Fig 18 (p33) Eigenvalues fit into clear areas of the

Tr-det plane. Eg. $\det(A) < 0 \Rightarrow$ saddles

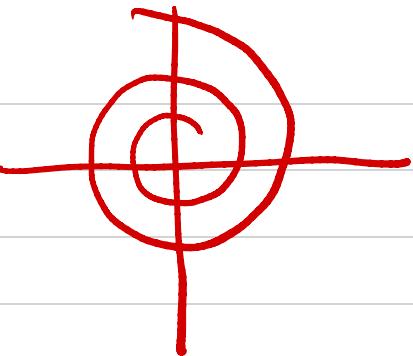
why because $\det(A) = \lambda_1 \lambda_2$ and $\lambda_1 \lambda_2 < 0$ only if
 λ_1, λ_2 are of opposite sign, etc.

WK8,9



nodes

$$J > 0$$

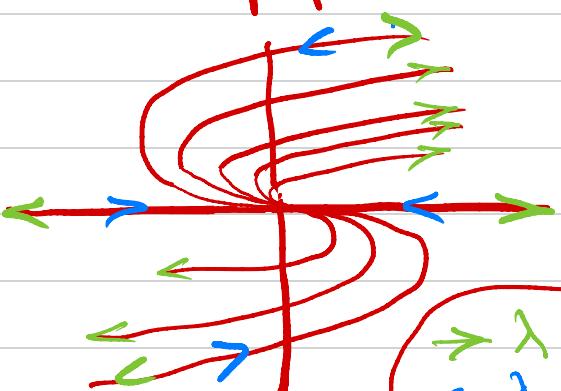


$$J < 0$$

improper node.

"Half way
between a node
and a spiral"

= IMPROPER
NODE

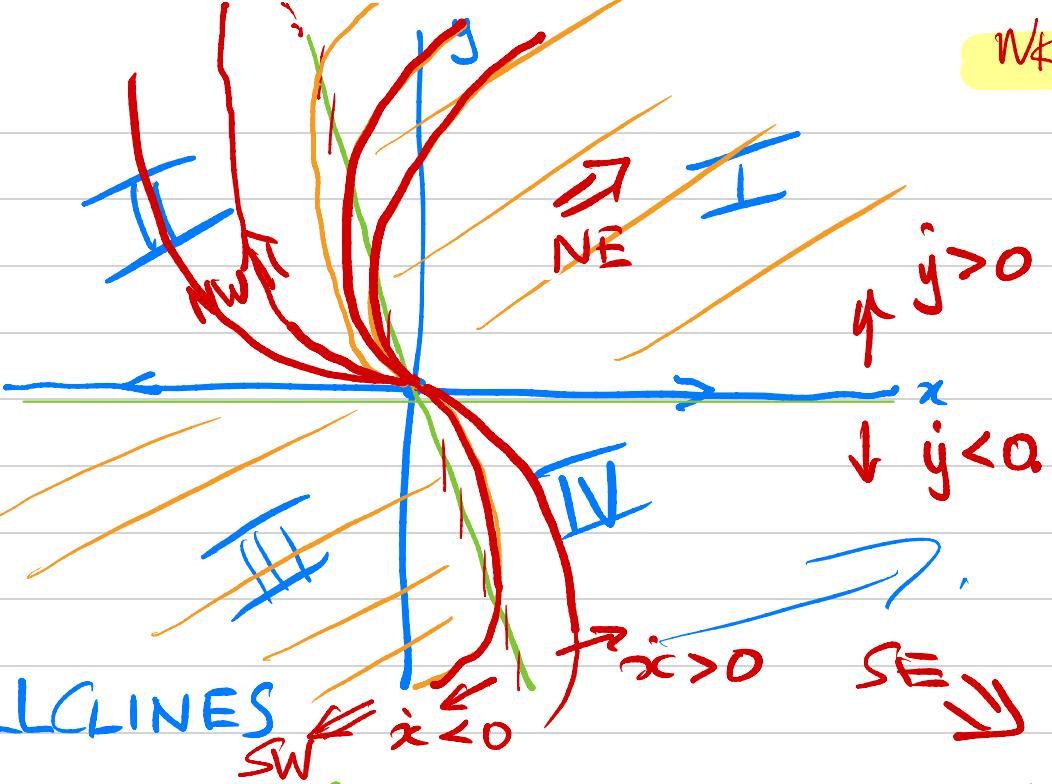


→ λ positive
← λ negative



$$A = \begin{bmatrix} 2 & 1 \\ 0 & 2 \end{bmatrix}$$

Wk 8.10



$$\dot{x} = 2x + y$$

$$\dot{y} = 2y$$

Consider

$$\dot{x} = 0$$

$$2x + y = 0$$

$$\dot{y} = 0$$

$$2y = 0, y = 0$$

having horizontally

(no movement in x -dirⁿ
 \therefore vertical dirⁿ only).

$$A = \begin{bmatrix} 4 & -1 \\ 1 & 2 \end{bmatrix}$$

$\text{Det}(\lambda I - A) = 0$ - the eigenvalue equation

WKB, 11

$$\text{Tr}(A) = 6$$

$$\lambda = \frac{6 \pm \sqrt{6^2 - 4 \cdot 9}}{2}$$

$$\text{Det}(A) = 9$$

Z

Reminder!

$$\lambda = 3 \text{ (rep).}$$

unstable improper node

Nullclines are pts where
 $\dot{x} = 0$ or $\dot{y} = 0$

A repeated root and non-diagonal

$$x = 0$$

$$4x - y = 0$$

↑?

$$y = 0$$

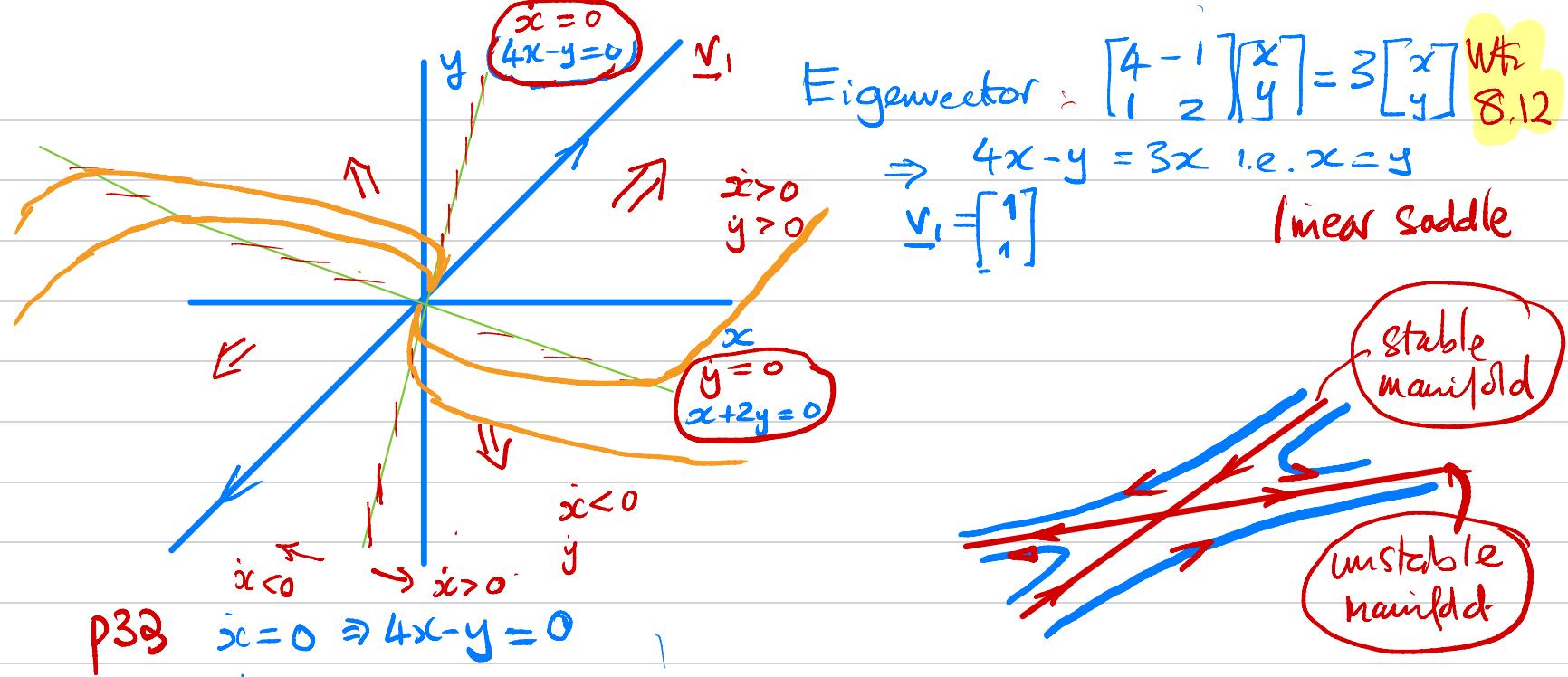
$$x + 2y = 0 \leftrightarrow ? \text{ Horizontally} \quad \dots$$

Any other isocline can be considered

e.g. $\frac{dy}{dx} = -1$

$$\frac{4x - y}{x + 2y} = -1 \Leftrightarrow 4x - y = -x - 2y$$

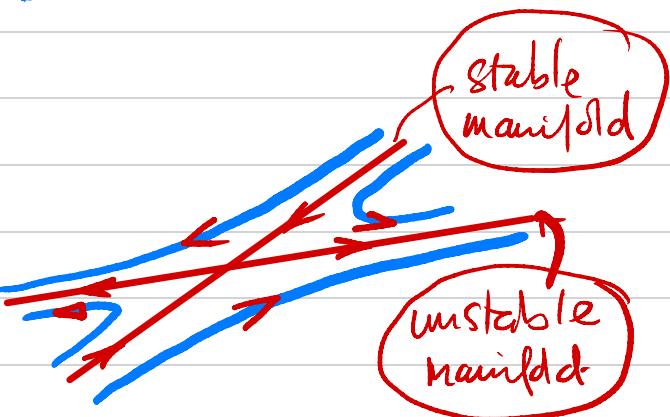
$$5x + y = 0$$



p33

$$\dot{x} = 0 \Rightarrow 4x - y = 0$$

$$\dot{y} = 0 \Rightarrow x + 2y = 0$$



Demos at end of Chapter 4
 "Attracting" is a loose definition
 and will not be used in the
 final exam.

There is much to come to terms with in today's class
But we do have practise in the last chapter of

- Locating fixed points
- investigating eigenvalues and eigenvectors
 - nullclines
 - general direction of solution curves (NE, etc)
 - sketching phase portraits.
 -